



# Neutrino Oscillation Parameters: Results and Prospects – II

Sandhya Choubey

Harish-Chandra Research Institute, Allahabad, India



III International Pontecorvo Neutrino Physics School  
September 16-26, 2007, Alushta, Ukraine



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- Why disappearance experiments?
- Because statistics there are very large



# The Unanswered Questions

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- Is there CP violation in the lepton sector?
  - Main channel to see  $\delta_{CP}$ 
    - $\nu_\mu \rightarrow \nu_e$  or  $\nu_e \rightarrow \nu_\mu$ :  $P_{\mu e}$  or  $P_{e\mu}$ ; Appearance Expts
    - Also possible using  $\nu_\mu \rightarrow \nu_\mu$   $P_{\mu\mu}$ ; Disapp Expts



# Measuring $\theta_{13}$ , $\delta_{CP}$ and $sgn(\Delta m_{31}^2)$



# Measuring $\theta_{13}$ with Reactors

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$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e) \approx 1 - \sin^2 2\theta_{13} \sin^2 \left( \frac{\Delta m_{31}^2 L}{4E} \right) + O(\alpha^2)$$



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- Proposals include:  
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- Sensitivity (90% C.L.):

$$\sin^2 2\theta_{13} < 0.032 \quad \text{Double Chooz}$$

$$\sin^2 2\theta_{13} < 0.009 \quad \text{Reactor II}$$



# The $\nu_e \rightarrow \nu_\mu$ “Golden” Channel

$$P_{app}^{vac} \simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin^2 \Delta$$

$$\pm (\alpha \Delta) \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \sin \delta_{CP} \sin \Delta \sin \Delta$$

$$+ (\alpha \Delta) \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta_{CP} \cos \Delta \sin \Delta$$

$$+ \cos^2 \theta_{23} \sin^2 2\theta_{12} \sin^2(\alpha \Delta)$$

where  $\Delta \equiv \frac{\Delta m_{31}^2 L}{4E}$ ,  $\alpha \equiv \frac{\Delta m_{21}^2}{\Delta m_{31}^2}$  and we have expanded the probability in  $\alpha$  and  $\theta_{13}$  keeping only lower order terms



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- First term gives  $\theta_{13}$
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- Third term gives CP conserving part of the prob



# Bi-Probability Plots

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# E•Probability Plots

$$\begin{aligned} P_{app}^{vac} = & \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin^2 \Delta \\ & \pm (\alpha \Delta) \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \sin \delta_{CP} \sin \Delta \sin \Delta \\ & + (\alpha \Delta) \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta_{CP} \cos \Delta \sin \Delta \\ & + \cos^2 \theta_{23} \sin^2 2\theta_{12} \sin^2(\alpha \Delta) \end{aligned}$$



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$$P_\nu = A \cos \delta_{CP} + B \sin \delta_{CP} + C$$
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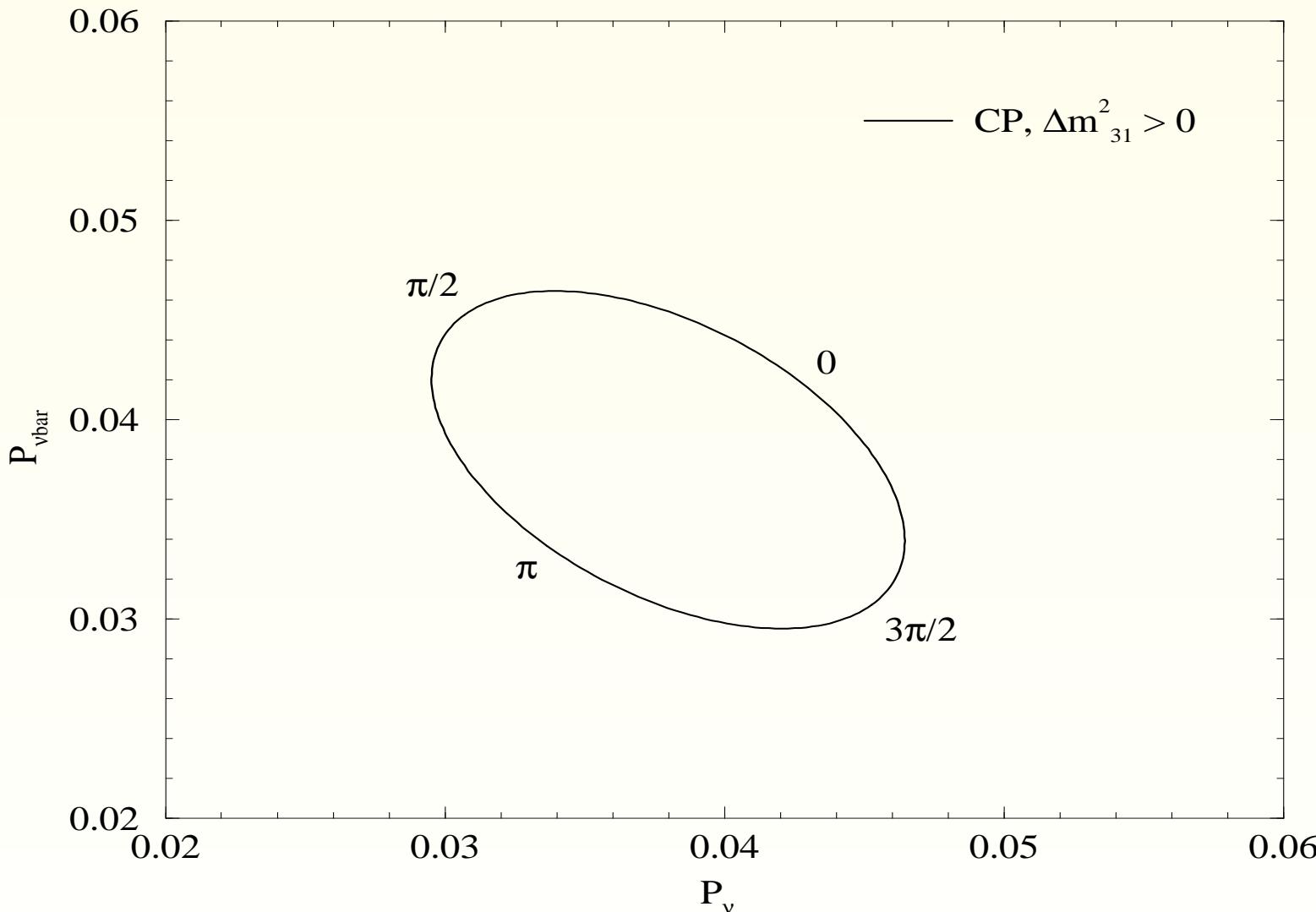
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- This is the eqn of an ellipse in the  $P_\nu - P_{\bar{\nu}}$  plane. These are called “**bi-probability plots**”. Major (minor) axes measure the amplitude of  $\sin \delta_{CP}$ ( $\cos \delta_{CP}$ )term

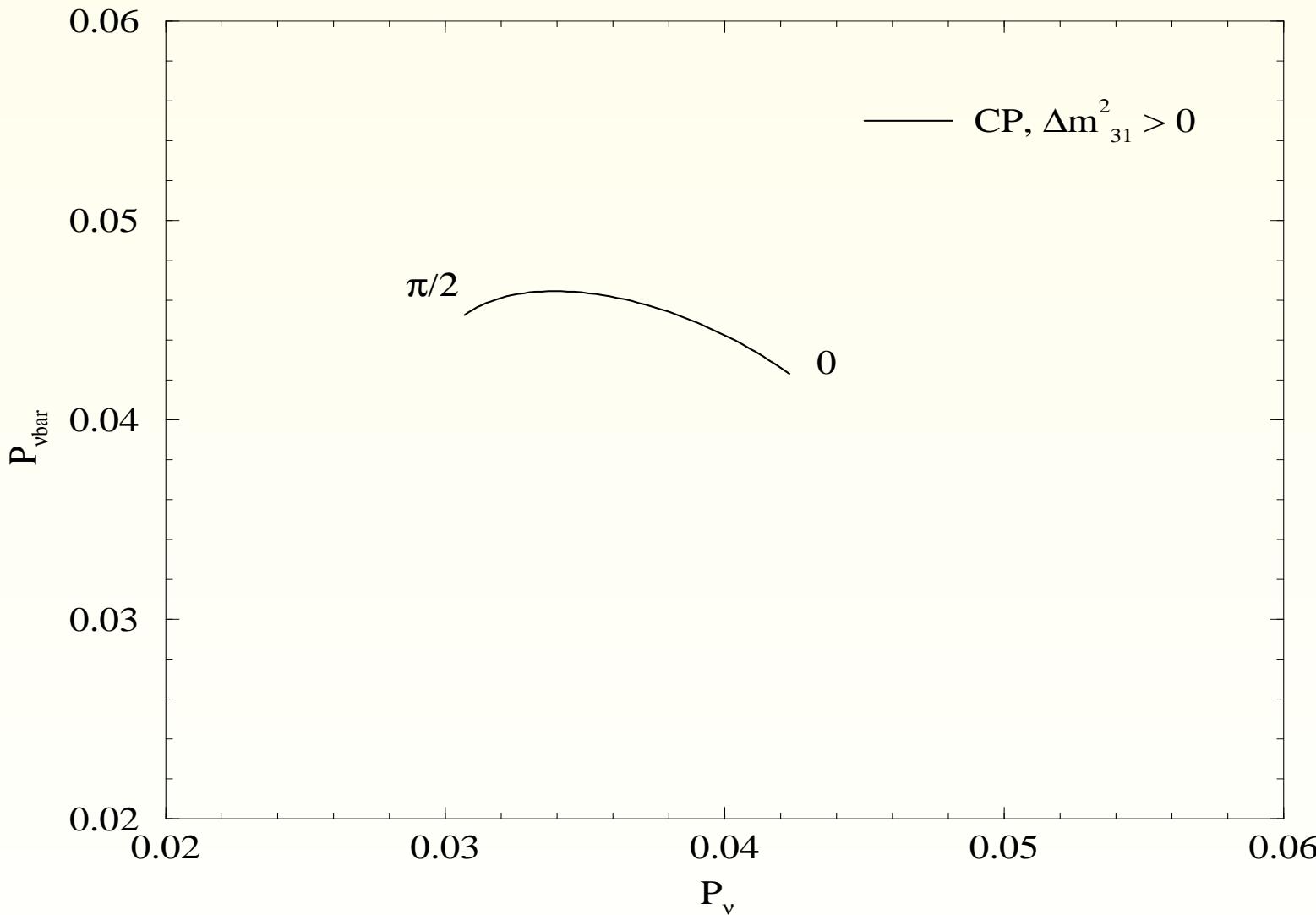


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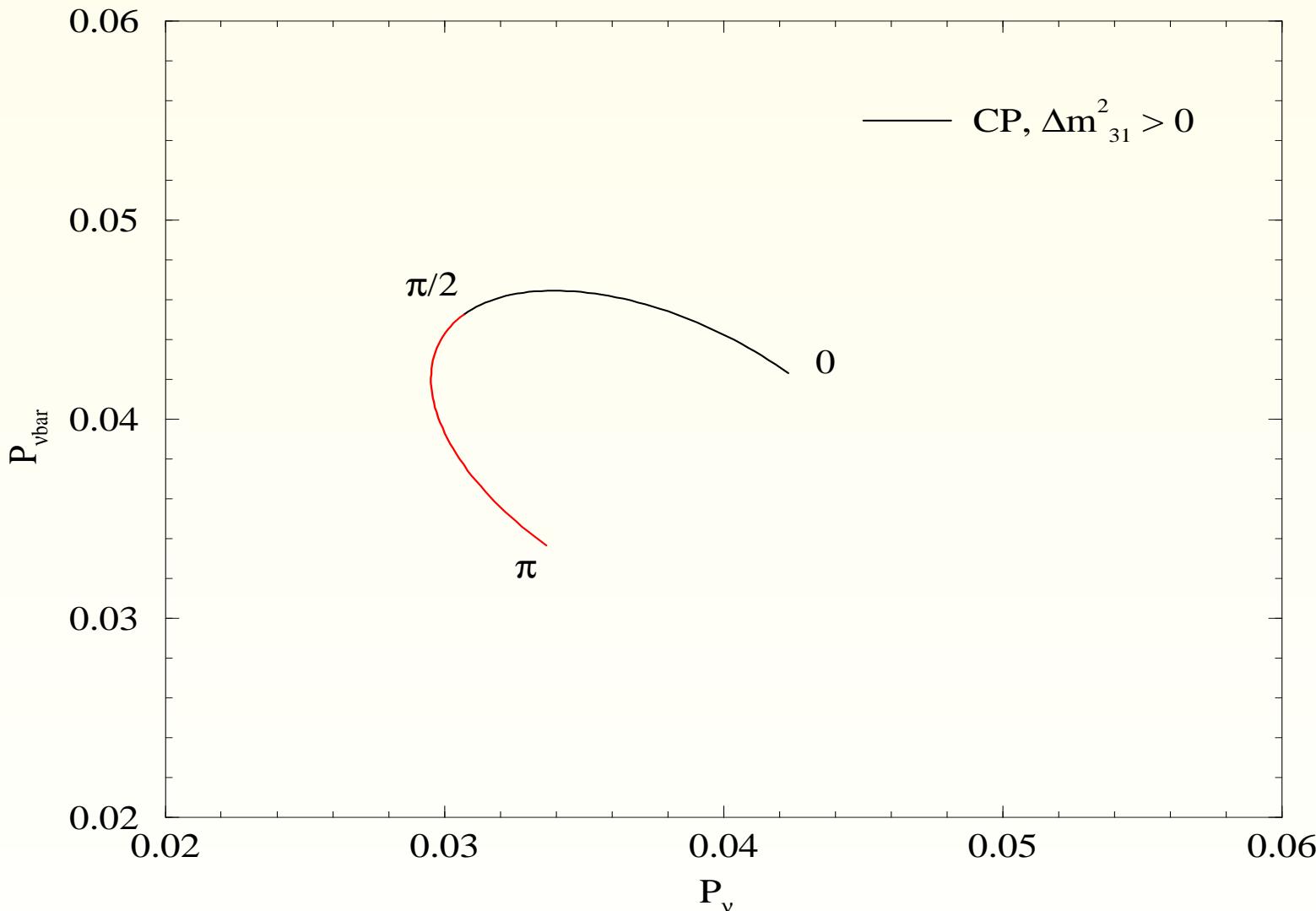


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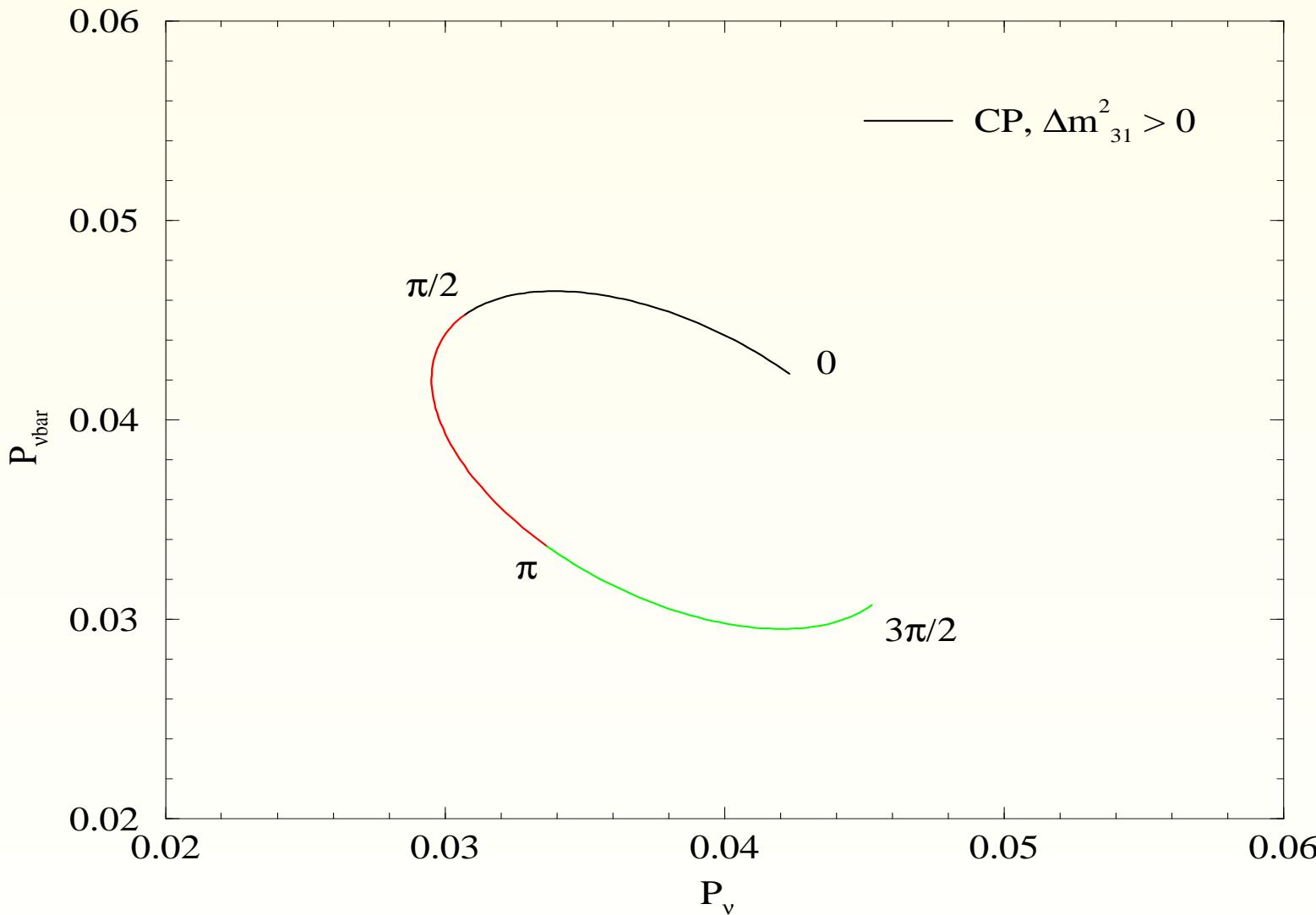


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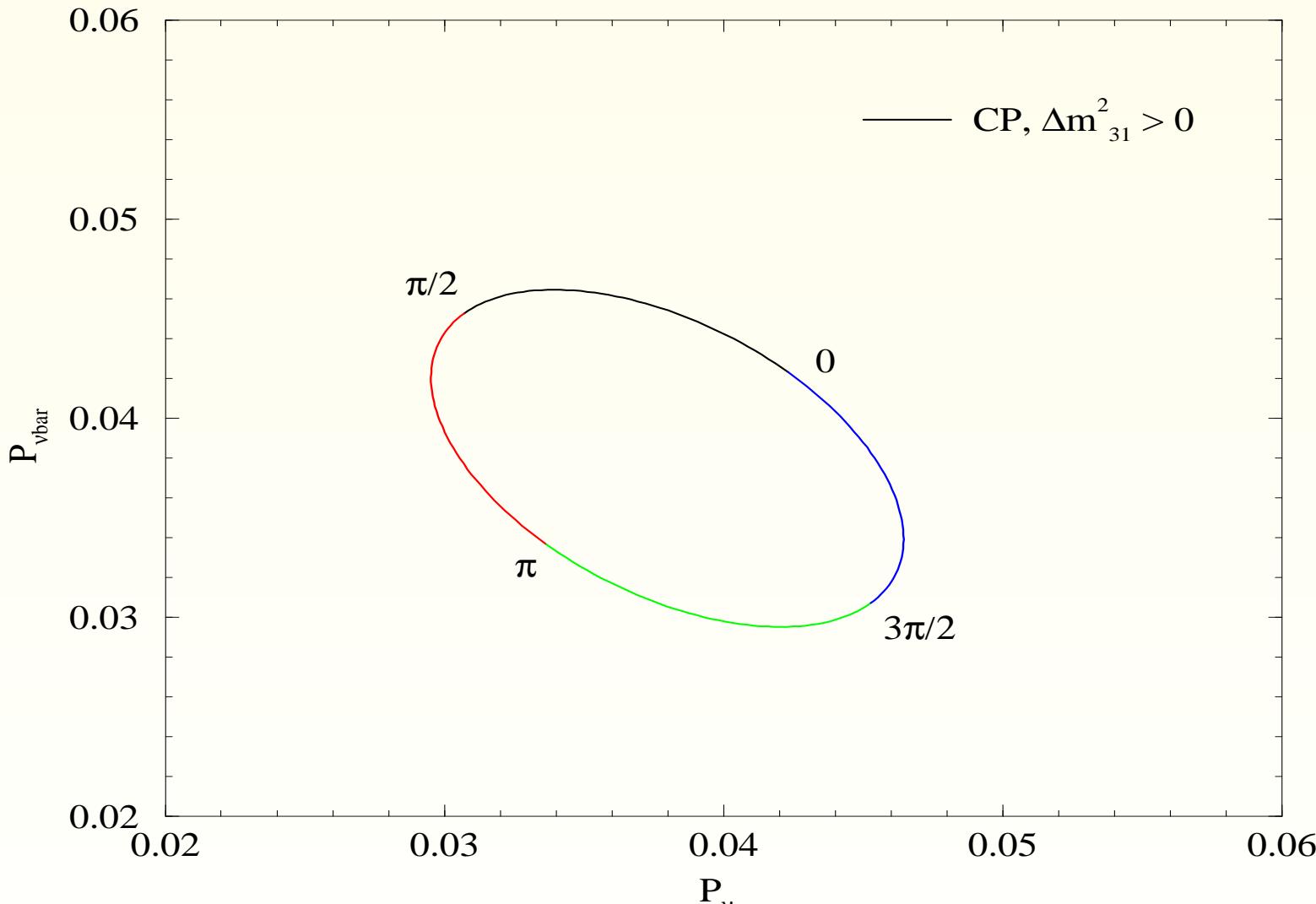


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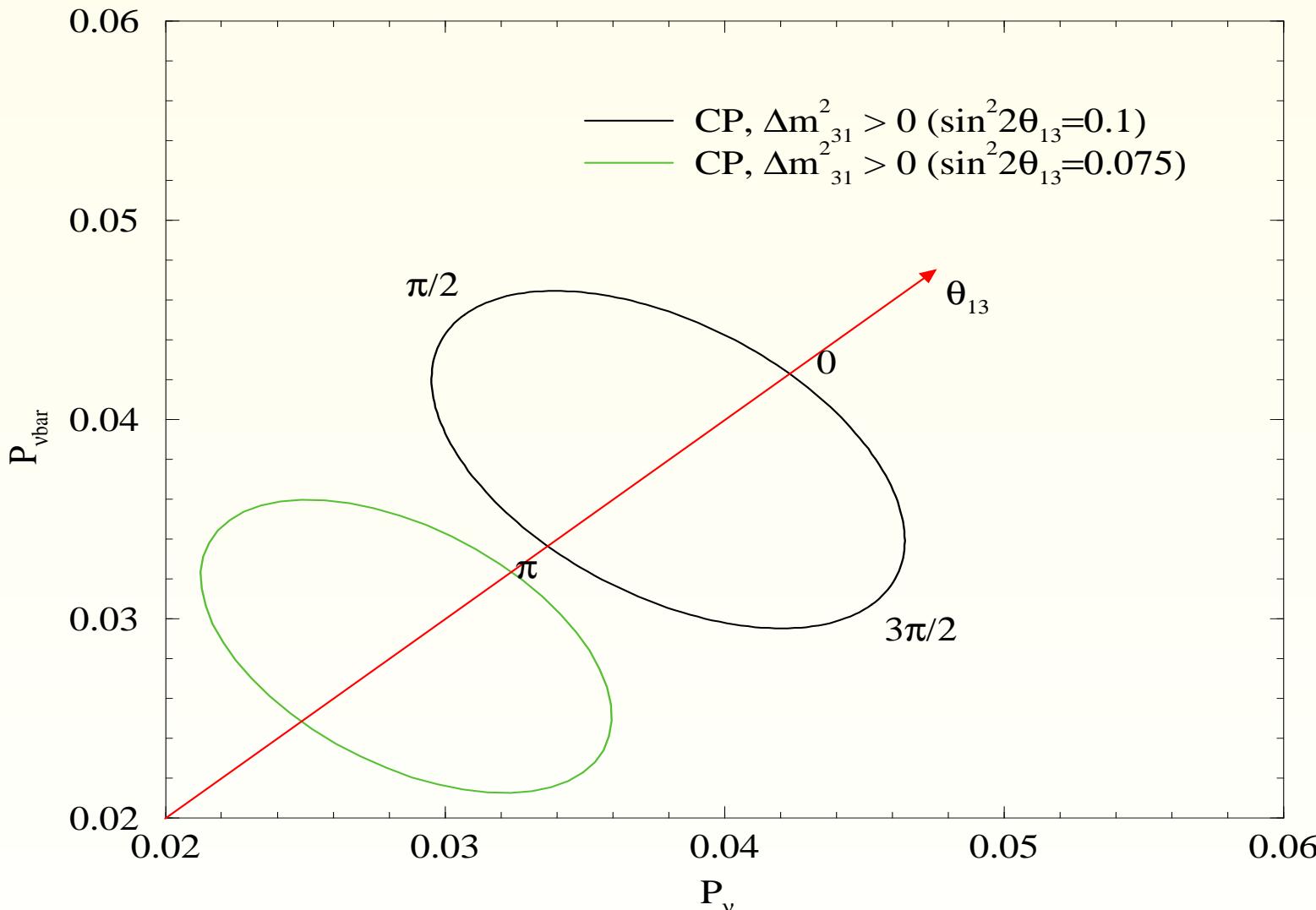
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- The ellipse moves anticlockwise for  $\Delta m^2_{31} > 0$ .
- You can check that it goes clockwise for  $\Delta m^2_{31} < 0$ .



# Bi-Probability Plots



Smaller  $\theta_{13} \Rightarrow \delta_{CP}$  measurement more difficult .



# Degeneracies

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# $Sgn(\Delta m_{31}^2)$ Degeneracy



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Do the simultaneous transformation

$$\delta_{CP} \rightarrow \pi - \delta_{CP}$$

$$\Delta m_{31}^2 \rightarrow -\Delta m_{31}^2$$



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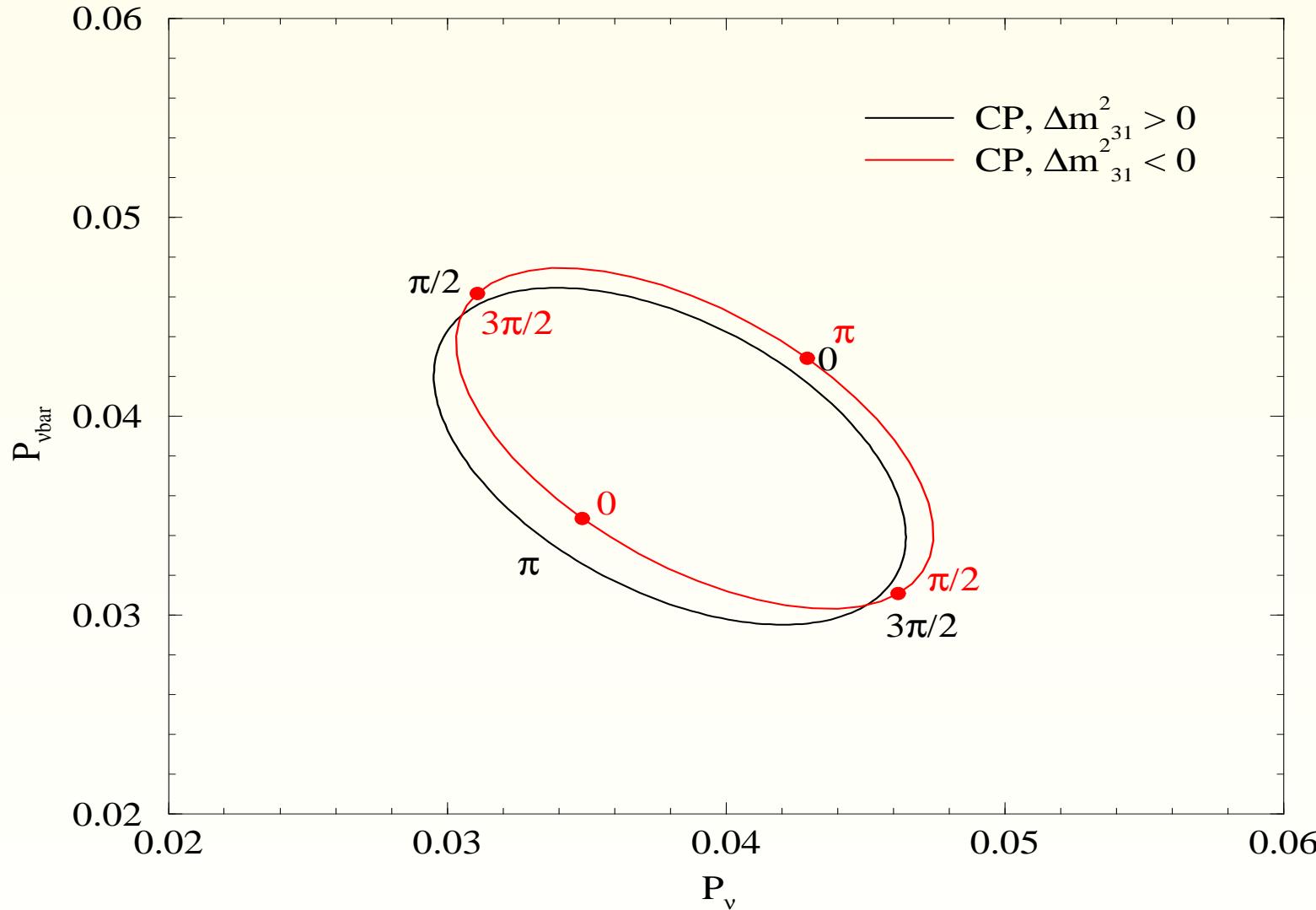
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- The expression for  $P_{\text{appearance}}$  remains invariant
- Since  $Sgn(\Delta m_{31}^2)$  is unknown, there will always be an ambiguity in the measured value of  $\delta_{CP}$



# $Sgn(\Delta m_{31}^2)$ Degeneracy





# Octant of $\theta_{23}$ Degeneracy



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- We only have measurement on  $\sin^2 2\theta_{23}$ .
- For every non-maximal  $\sin^2 2\theta_{23}$ , there are 2 possible  $\sin^2 \theta_{23}$



# Octant of $\theta_{23}$ Degeneracy



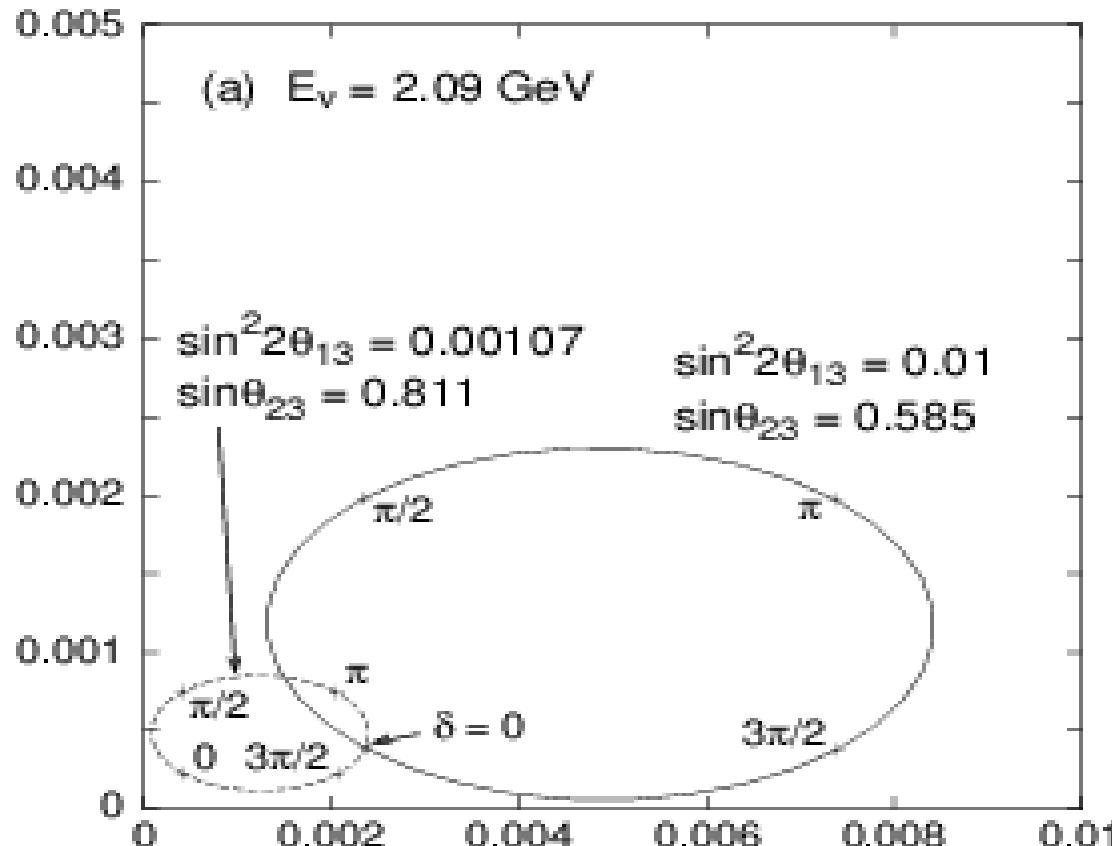
$$\begin{aligned} P_{\text{appearance}} &= \sin^2 \theta_{23} \sin^2 2\theta_{13} \sin^2 \Delta \\ &\pm (\alpha \Delta) \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \sin \delta_{CP} \sin \Delta \sin \Delta \\ &+ (\alpha \Delta) \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta_{CP} \cos \Delta \sin \Delta \\ &+ \cos^2 \theta_{23} \sin^2 2\theta_{12} \sin^2(\alpha \Delta) \end{aligned}$$

- We only have measurement on  $\sin^2 2\theta_{23}$ .
- For every non-maximal  $\sin^2 2\theta_{23}$ , there are 2 possible  $\sin^2 \theta_{23}$
- This will give 2 disjoint fitted value for  $\theta_{13}$ .



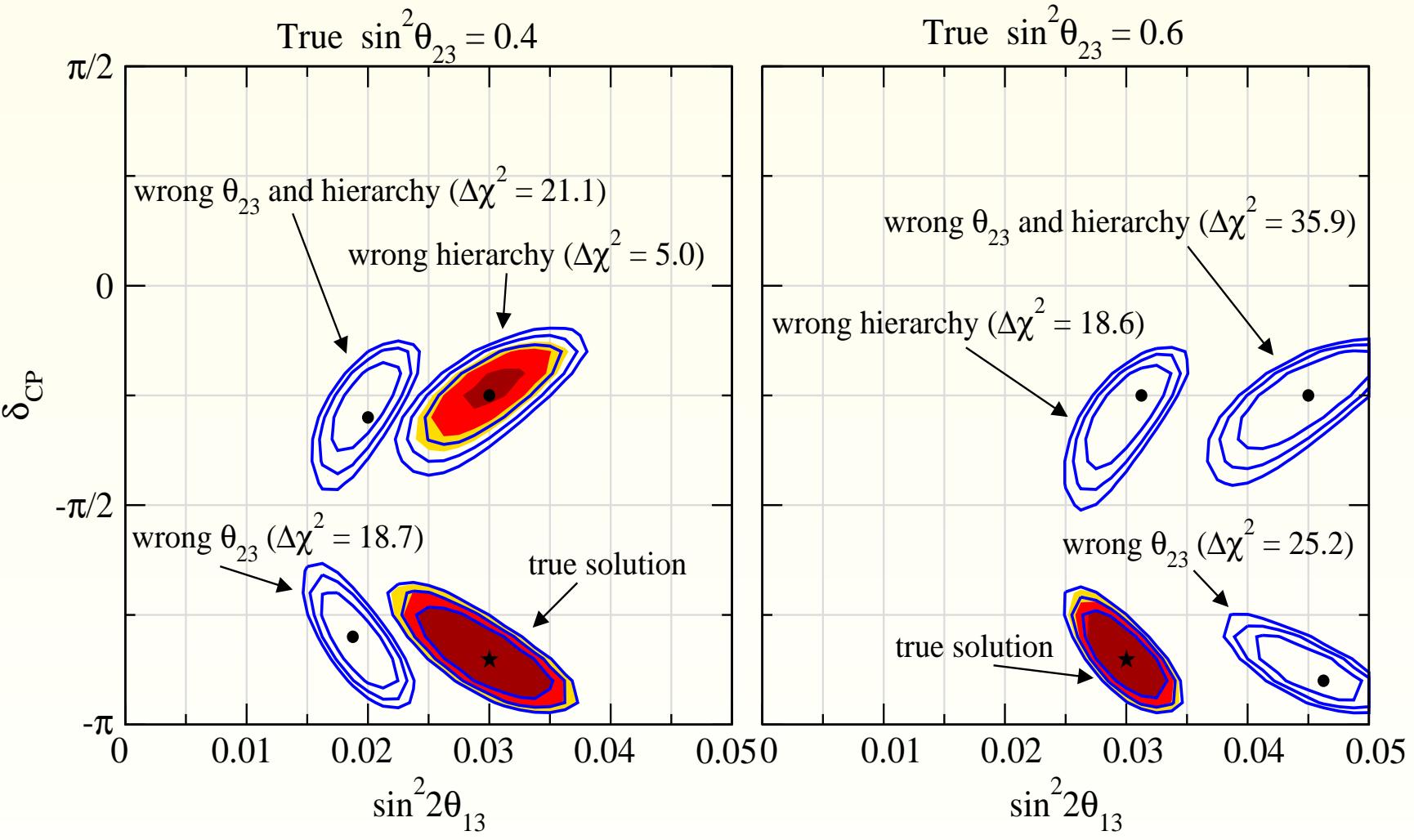
# Octant of $\theta_{23}$ Degeneracy

$L = 1290 \text{ km}, \delta m_{31}^2 = 3 \times 10^{-3} \text{ eV}^2, \delta m_{21}^2 = 5 \times 10^{-5} \text{ eV}^2$



Barger et al, hep-ph/0112119

- If  $\sin^2 2\theta_{23} = 0.9$  then both  $\sin^2 2\theta_{13} = 0.01$  and  $0.0011$  would fit the data from a given LBL expt.



Huber et al, hep-ph/0501037



# Intrinsic $(\delta_{CP}, \theta_{13})$ Degeneracy

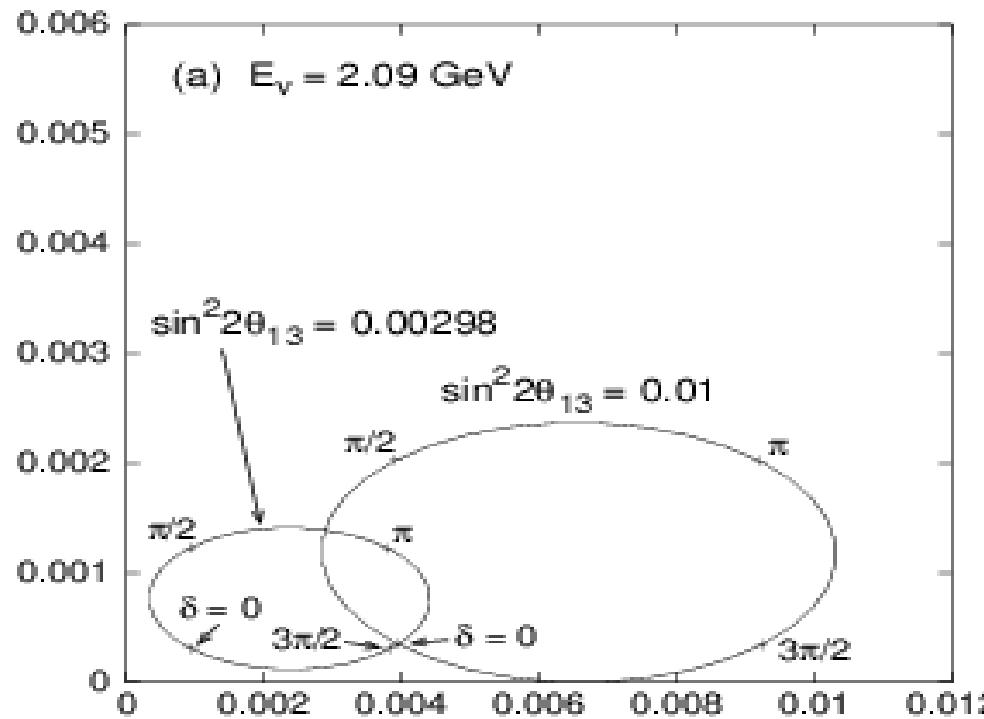
- For a fixed value of  $E$ , it might happen that

$$P_{app}(\theta_{13}, \delta_{CP}) = P_{app}(\theta'_{13}, \delta'_{CP})$$

$$\bar{P}_{app}(\theta_{13}, \delta_{CP}) = \bar{P}_{app}(\theta'_{13}, \delta'_{CP})$$

- $(\theta_{13}, \delta_{CP})$  (true solution) &  $(\theta'_{13}, \delta'_{CP})$  (fake solution)

$L = 1290$  km,  $\delta m^2_{31} = 3 \times 10^{-3}$  eV $^2$ ,  $\delta m^2_{21} = 5 \times 10^{-5}$  eV $^2$





Up to Eight-Fold Degeneracy Expected



# The $\nu_e \rightarrow \nu_\mu$ Channel in Matter



# The $\nu_e \rightarrow \nu_\mu$ Channel in Matter

$$P_{app} \simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \frac{\sin^2[(1 - \hat{A})\Delta]}{(1 - \hat{A})^2}$$

$$\pm \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \sin \delta_{CP} \sin(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 - \hat{A})\Delta]}{(1 - \hat{A})}$$

$$+ \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta_{CP} \cos \Delta \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 - \hat{A})\Delta]}{(1 - \hat{A})}$$

$$+ \alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}$$

- $\hat{A} = \frac{\pm A}{\Delta m_{31}^2}$

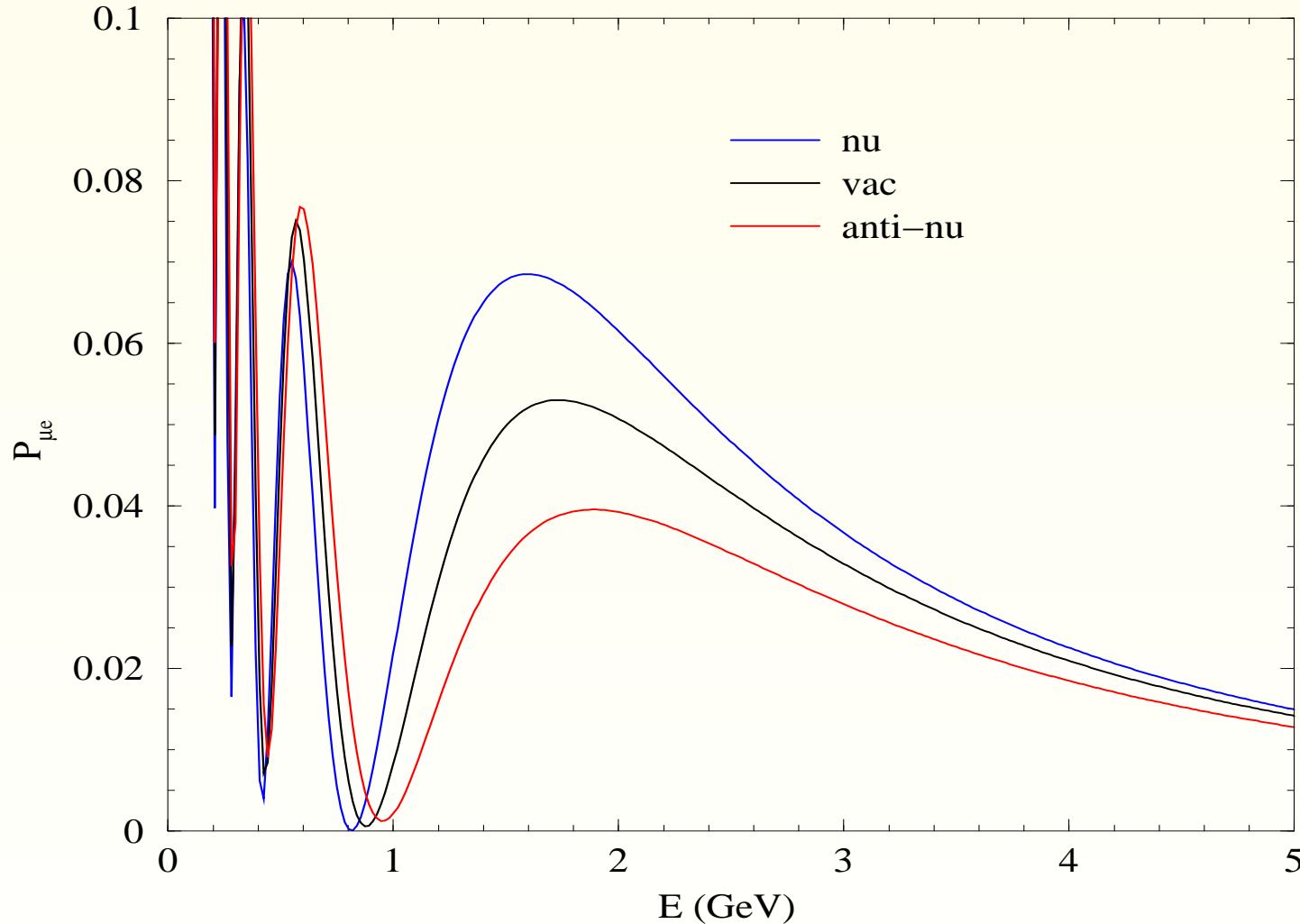
- $A$  is positive for neutrinos

- $A$  is negative for antineutrinos



# The $\nu_e \rightarrow \nu_\mu$ Channel in Matter

$L=1000\text{km}$ ,  $\Delta m^2_{31}=0.002$ ,  $\Delta m^2_{21}=8\times 10^{-5}$ ,  $s^2_{23}=0.5$ ,  $s^2_{12}=0.31$ ,  $\sin^2\theta_{13}=0.1$ ,  $\delta=0$



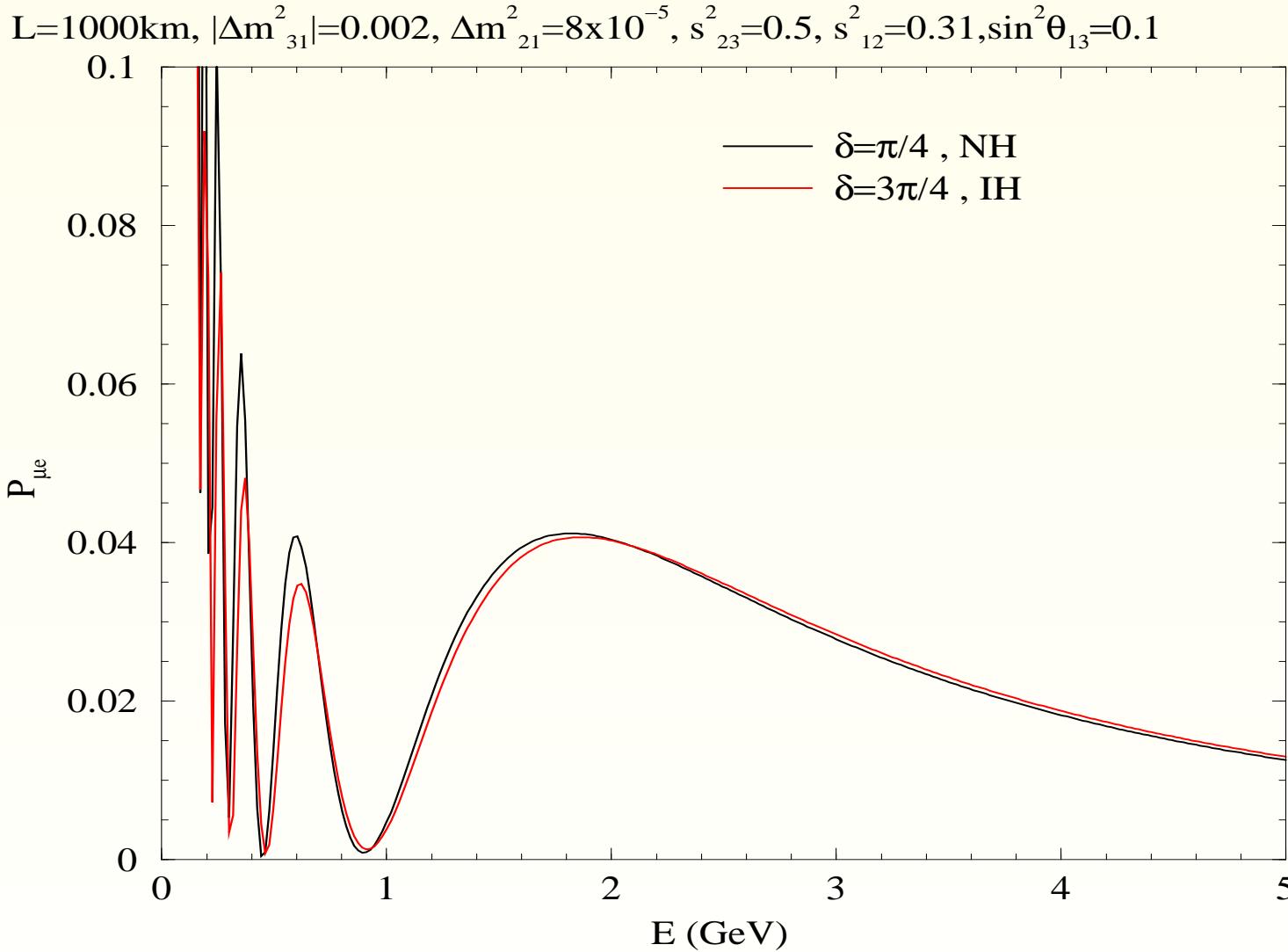


# Resolving the Eight-Fold Degeneracy

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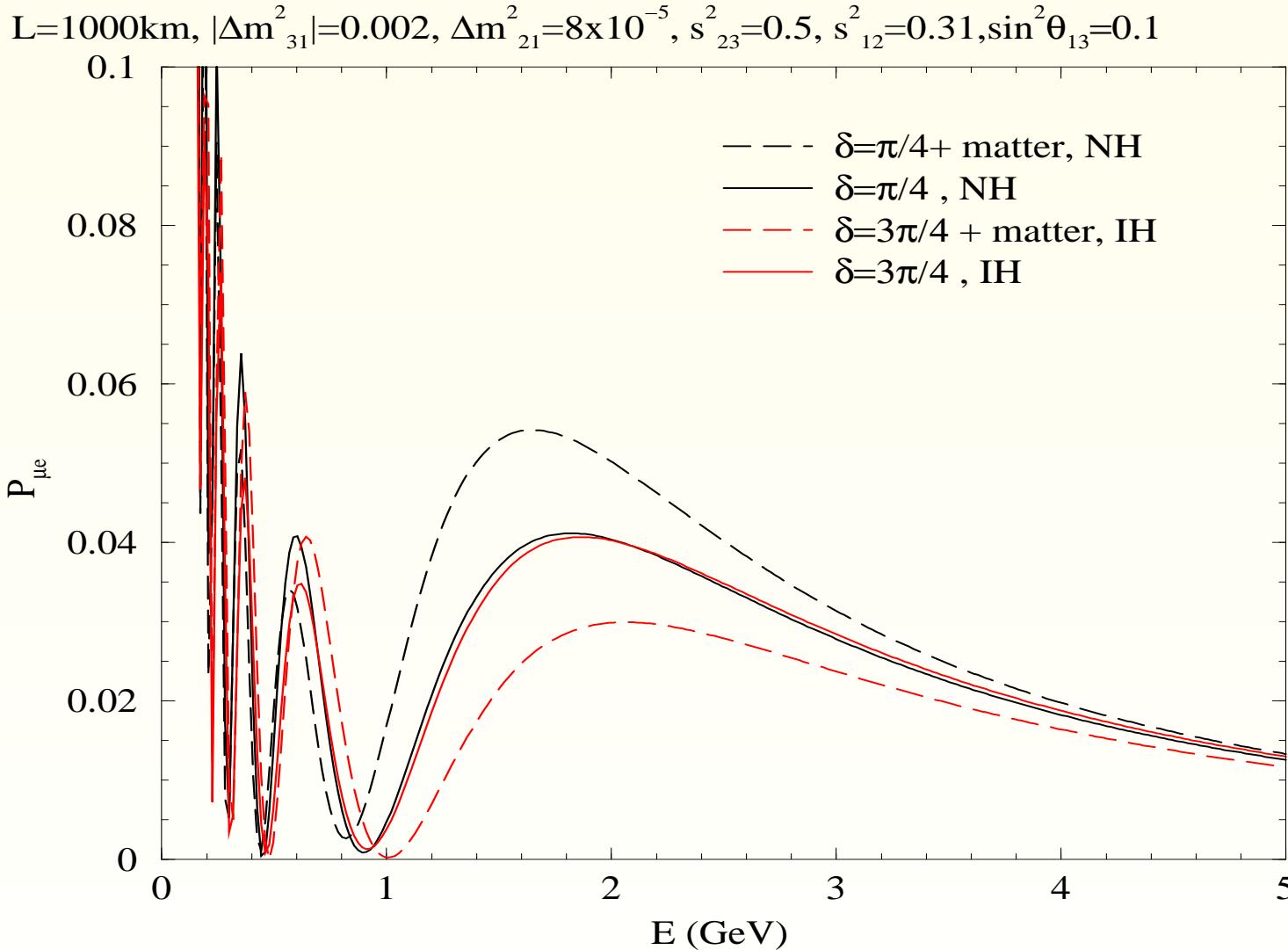


# Resolving the $Sgn(\Delta m_{31}^2)$ Degeneracy



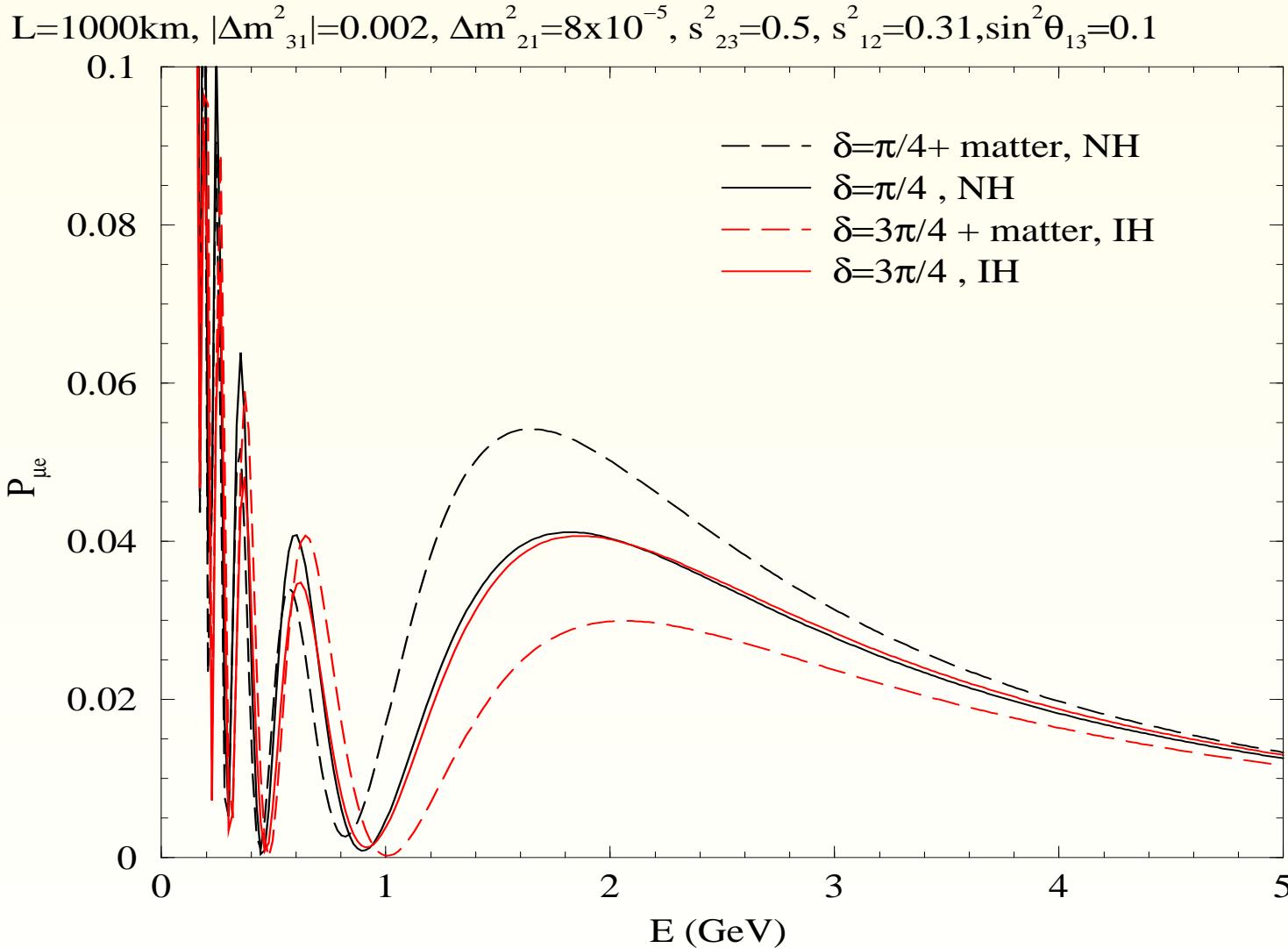


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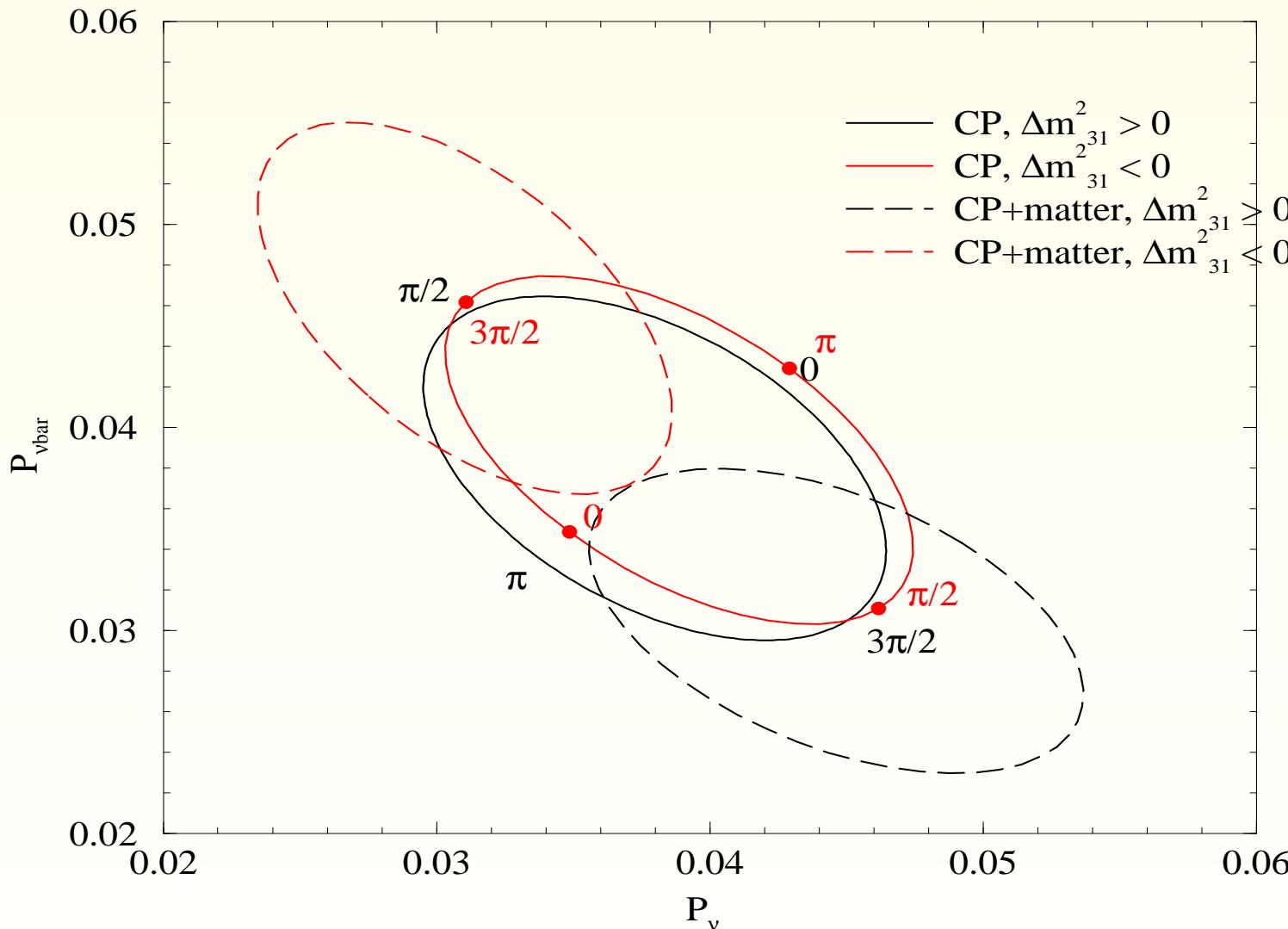
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- Matter effects break the  $Sgn(\Delta m_{31}^2)$  Degeneracy



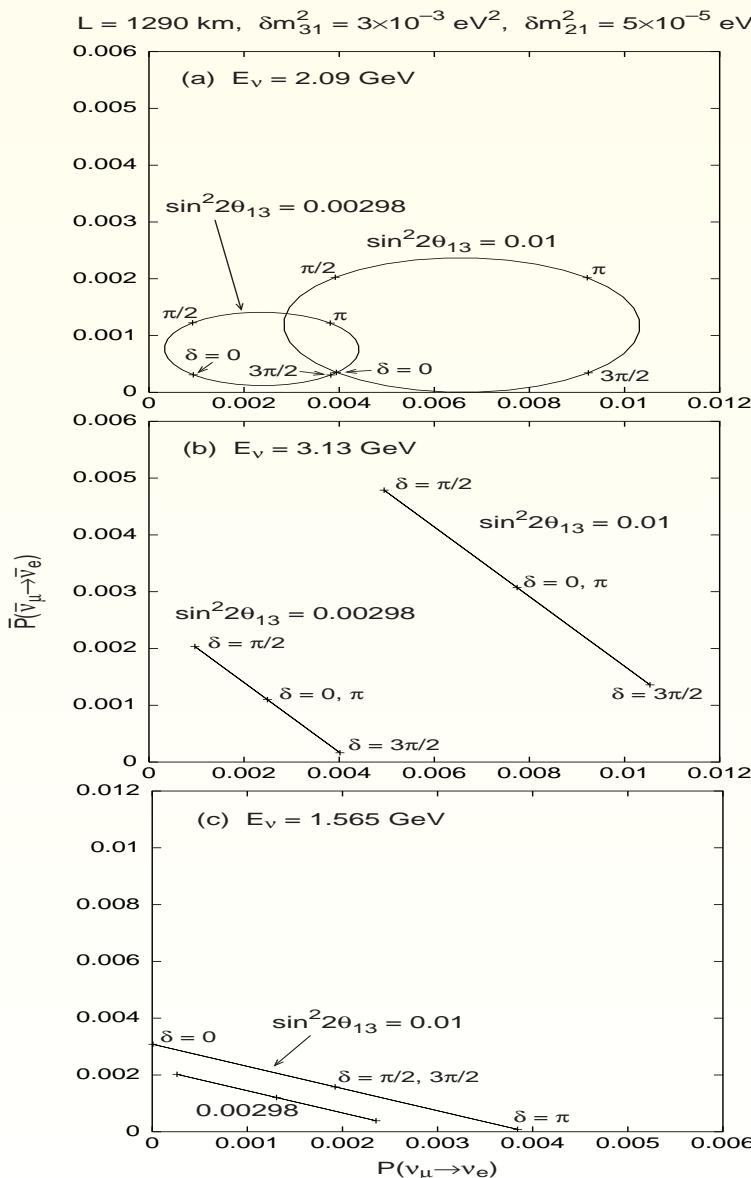
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- Matter effects break the  $Sgn(\Delta m_{31}^2)$  Degeneracy



# Resolving the Intrinsic $(\delta_{CP}, \theta_{13})$ Degeneracy



- Choose  $\Delta = m\pi/2$ .
- The  $\cos \delta_{CP}$  term vanishes for  $\Delta = (n - \frac{1}{2})\pi$ .
- The  $\sin \delta_{CP}$  term vanishes for  $\Delta = n\pi$ .
- Ellipse collapse to a line
- Ambiguity resolved.
- Better to work with  $\Delta = (n - \frac{1}{2})\pi$ .
- That is where we will directly see CPV.
- That's where we get the oscillation maxima.

Barger et al, hep-ph/0112119



# Resolving the Octant of $\theta_{23}$ Degeneracy



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- Using atmospheric neutrino data in
  - ✓ Megaton water detectors:
    - ★  $\Delta m_{21}^2$  driven osc effects in sub-GeV electrons
    - ★  $\theta_{13}$  driven matter effects in multi-GeV electrons
  - ✓ Large magnetized iron detectors:
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- Using data from next generation reactor neutrinos
- Other ways using LBL expts have been suggested.



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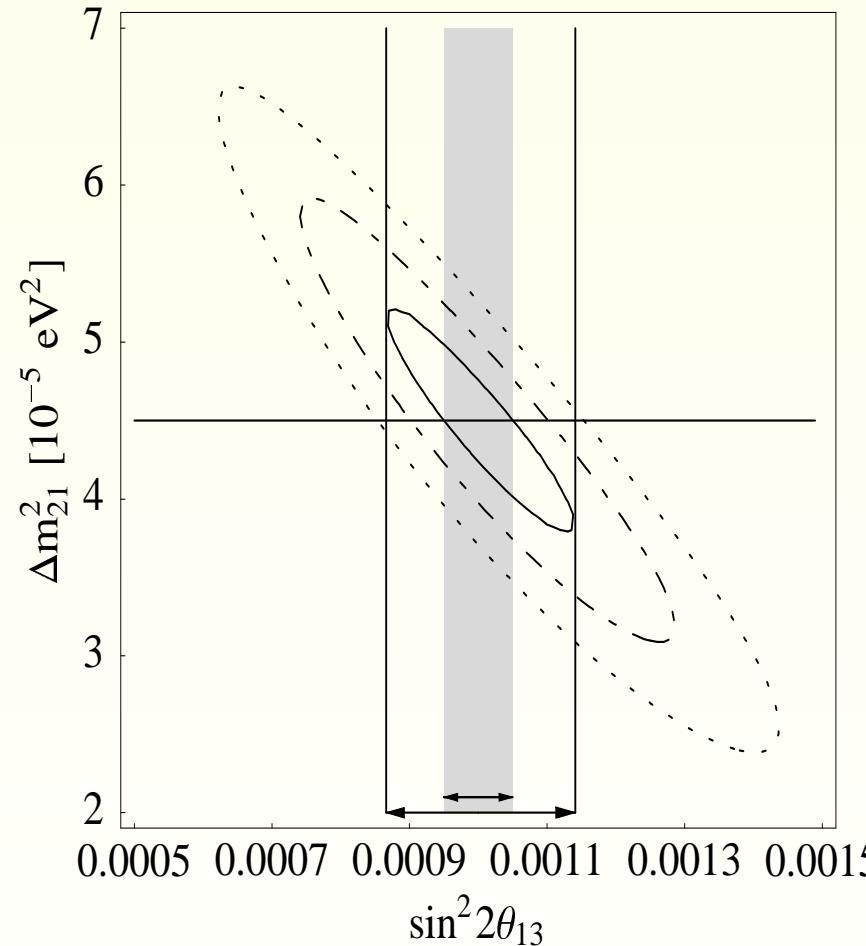
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- This correlation between parameters leads to increase in the error in the measured value of the individual parameters.



# Correlations

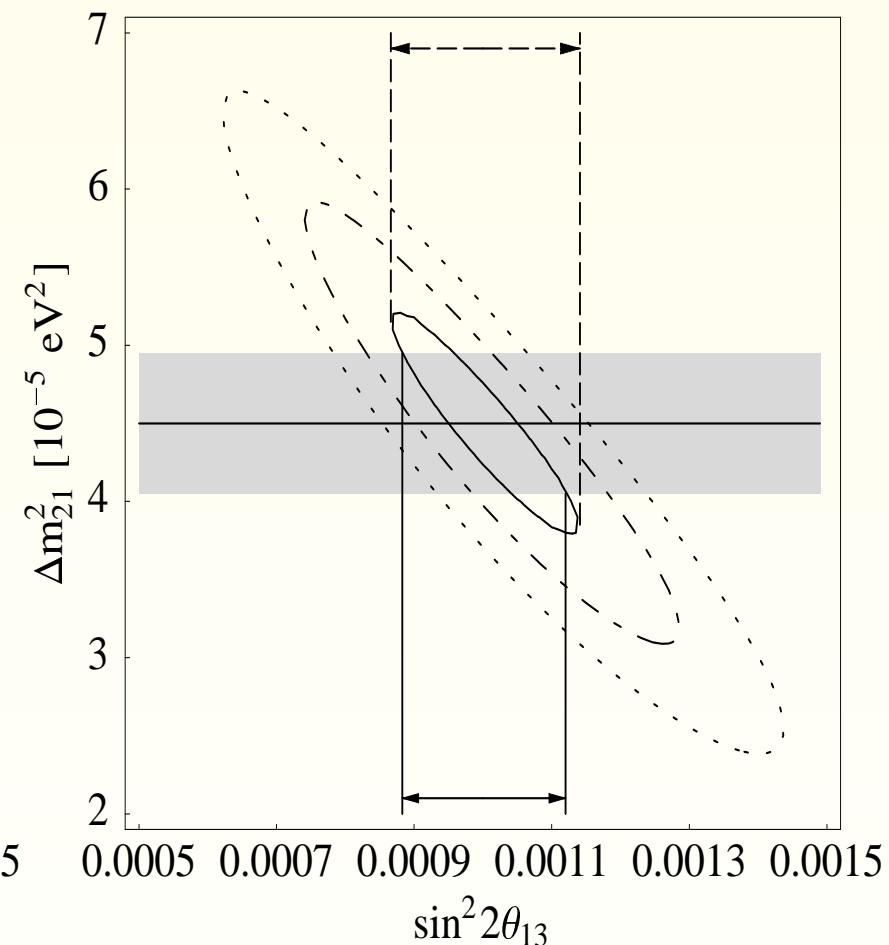
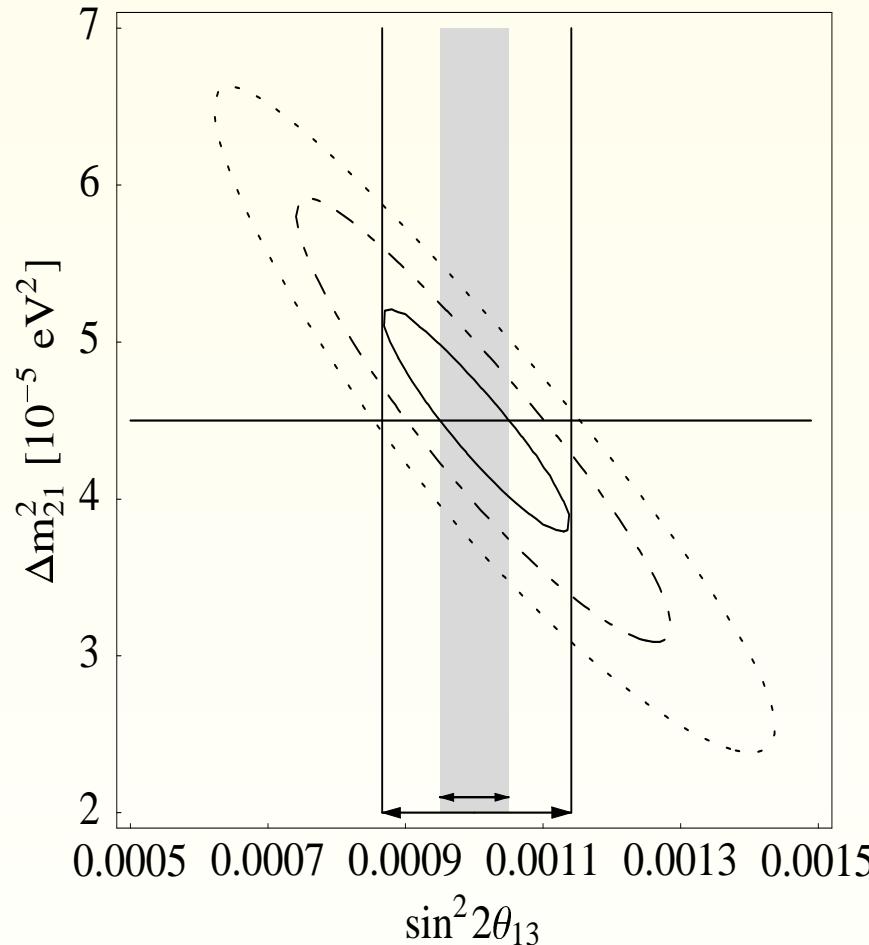


Huber et al, hep-ph/0204352

- Note how correlations are increasing the spread.



# Correlations



- Input from an external measurement helps. [Huber et al, hep-ph/0204352](#)
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# Differentiating Correlations and Degeneracies

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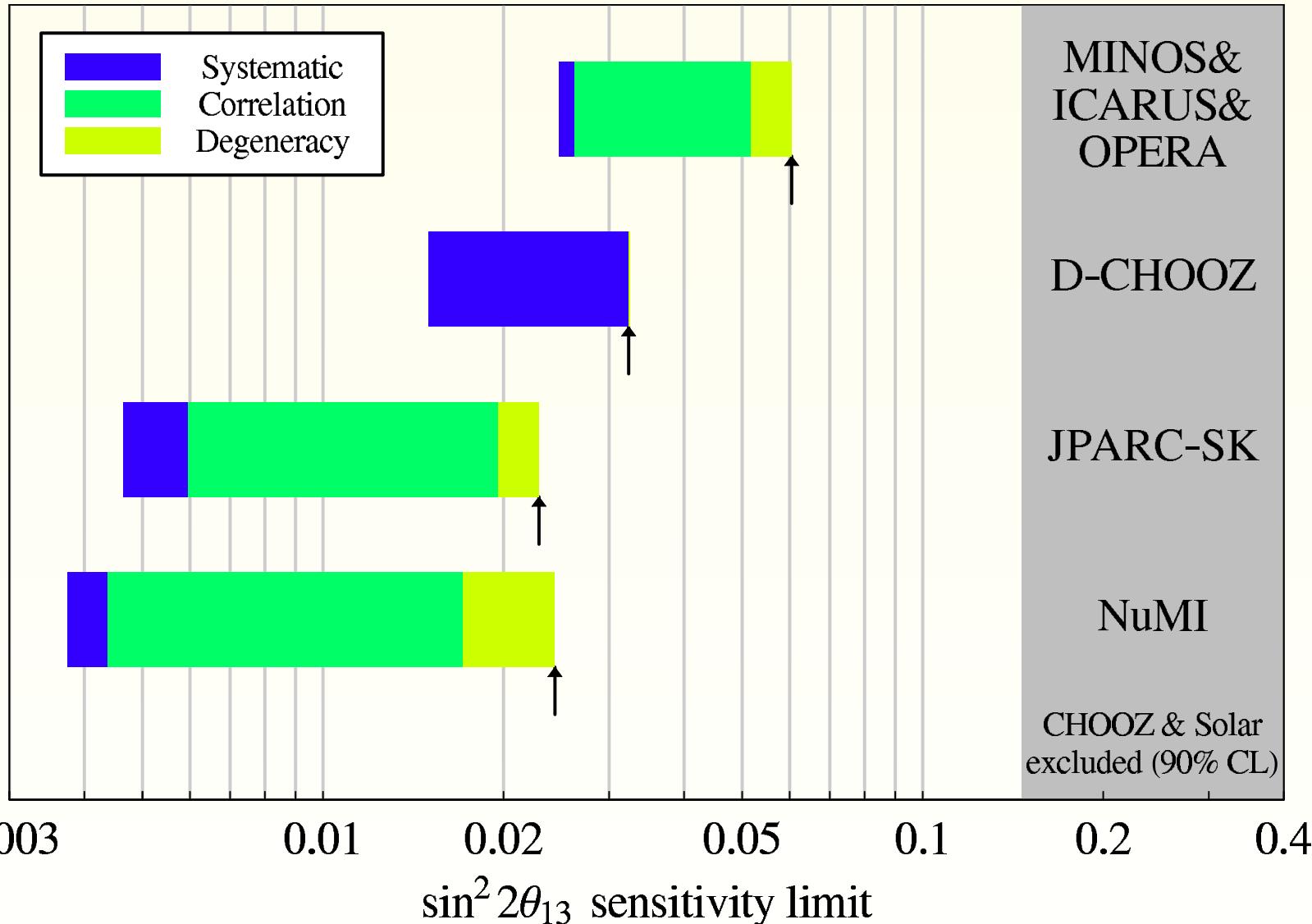
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- The only way to reduce impact of correlations is to combine experiments with different characteristics.



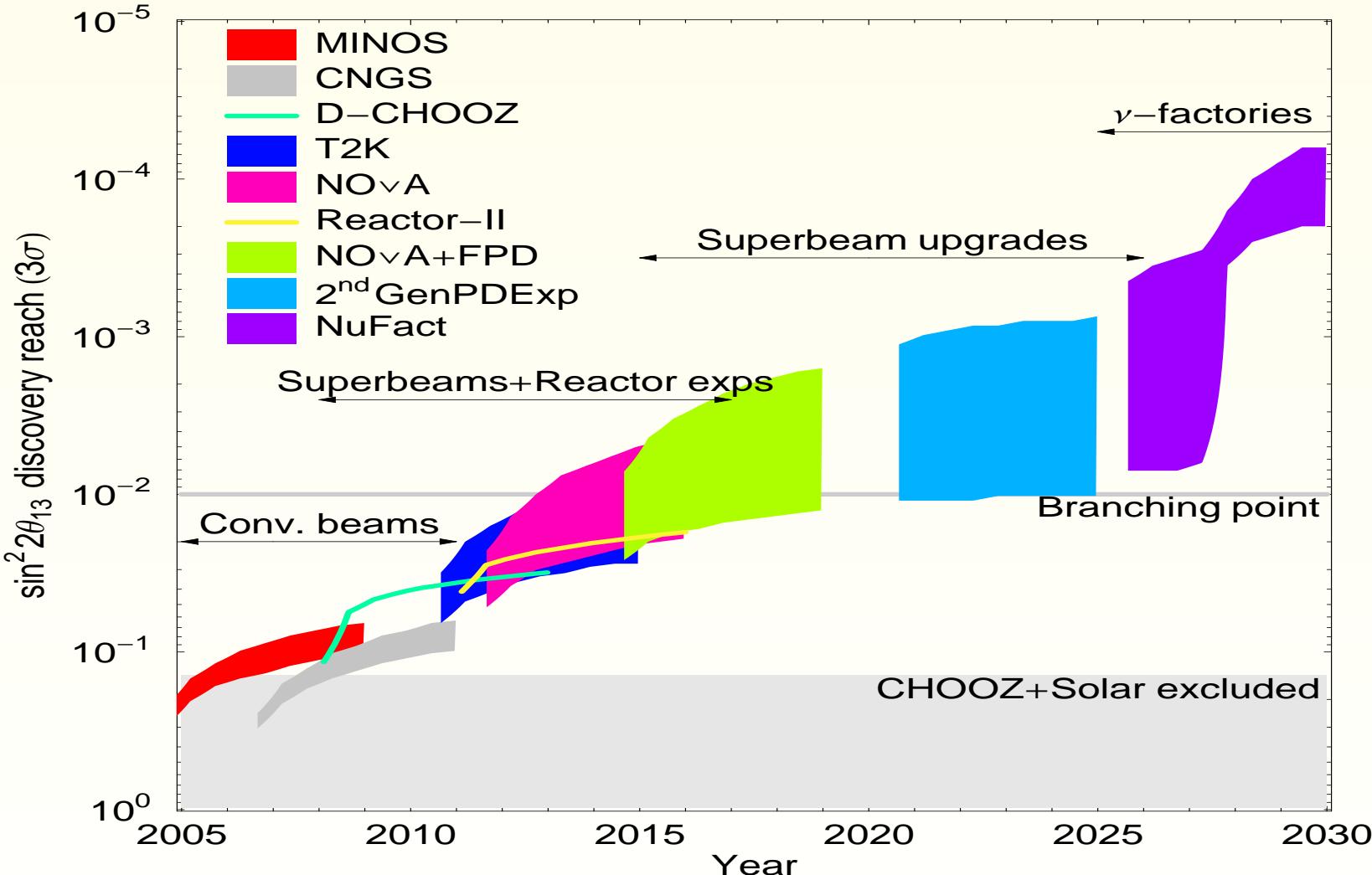
# Sensitivity of Near-Future Experiments to $\theta_{13}$



Huber *et al.*, hep-ph/0403068



# Sensitivity of Near-Future Experiments to $\theta_{13}$



Albrow *et al.*, hep-ex/0509019

# Sensitivity of Far-Future LBL Experiments

- # ● SuperBeam $\Rightarrow P_{\mu e}$



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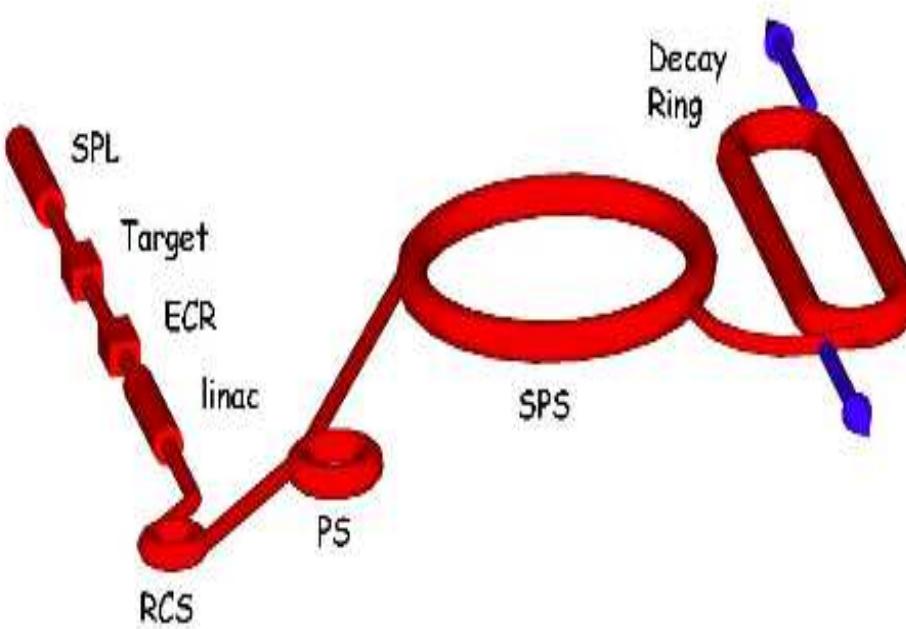
- SuperBeam  $\Rightarrow P_{\mu e}$
- Beta-Beam  $\Rightarrow P_{e\mu}$
- Neutrino Factory  $\Rightarrow P_{e\mu}$



# BetaBeams

- Beta-Beams are produced from beta decay of accelerated radioactive ions, circulating in a storage ring

$$(A, Z) \rightarrow (A, Z + 1) + e^+ + \nu_e \quad \text{or} \quad (A, Z) \rightarrow (A, Z - 1) + e^- + \bar{\nu}_e$$



- Proton Driver – SPL ( $\approx 4$  GeV)
- Target
- Ion Source – Pulsed ECR
- Accelerators – linac, RCS, PS, SPS
- Storage Ring – 7000m; 2500m straight

Zucchelli, PLB 532, 166, (2002)



# BetaBeams

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- Pure beam with just one flavor



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---

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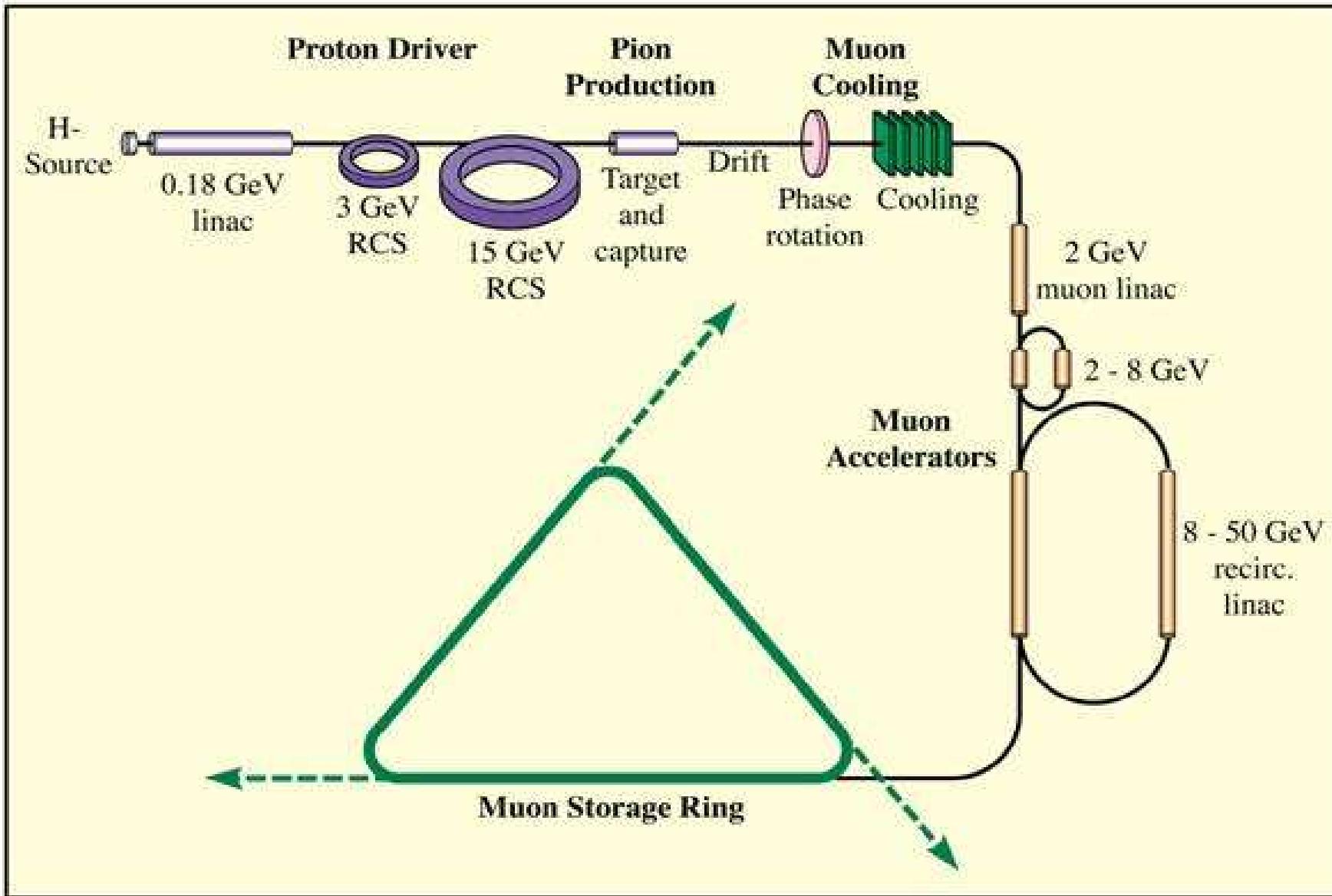
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- Can produce either  $\nu_e$  OR  $\bar{\nu}_e$  FLUX



# Neutrino Factory



Neutrino Factory at RAL



# Neutrino Factory

- Produced by decay of accelerated muons circulating in a storage ring:

$$\mu^+ \rightarrow e^+ + \nu_e + \bar{\nu}_\mu \text{ and/or } \mu^- \rightarrow e^- + \bar{\nu}_e + \nu_\mu$$



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- For a given muon sign in the ring, we will have muons of both signs in the detector. For a  $\mu^+$  source  $\bar{\nu}_\mu$  in the original beam will give rise to  $\mu^+$  in the detector (**right sign muons**), while  $\nu_e \rightarrow \nu_\mu$  oscillations will give  $\mu^-$  (**wrong sign muons**).



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- Charge ID is a MUST.
- Need magnetised detectors



# Sensitivity of Far-Future LBL Experiments

- SuperBeam  $\Rightarrow P_{\mu e}$
- Beta-Beam  $\Rightarrow P_{e \mu}$
- Neutrino Factory  $\Rightarrow P_{e \mu}$
- One has to find ways to kill the clone solutions
- Various ways have been suggested in the literature for this



# Killing the Clones

---

- Combining data from appearance experiments at different  $L$  and/or different  $E$ :

Barger, Marfatia, Whisnant, hep-ph/0206038

Barger, Marfatia, Whisnant, hep-ph/0210428

Burguet-Castell, Gavela, Gomez-Cadenas, Hernandez, Mena, hep-ph/0103258

Huber, Lindner, Winter, hep-ph/0211300

Mena and Parke hep-ph/0408070

Mena, Palomares-Ruiz, Pascoli, hep-ph/0504015

Mena, Palomares-Ruiz, Pascoli, hep-ph/0510182

Mena, Nunokawa, Parke, hep-ph/0609011

Minakata, Nunokawa, hep-ph/9706281

Minakata, Nunokawa, Parke, hep-ph/0301210

Ishitsuka, Kajita, Minakata, Nunokawa, hep-ph/0504026

Hagiwara, Okamura, Senda, hep-ph/0607255



# Killing the Clones

---

- Combining data from different channels:

The Silver Channel  $P_{e\tau}$

Autiero *et al.*, hep-ph/0305185

Donini, Meloni, Migliozzi, hep-ph/0206034

Disappearance Channel  $P_{\mu\mu}$

Donini, Fernandez-Martinez, Meloni, Rigolin, hep-ph/0512038

Donini, Fernandez-Martinez, Rigolin, hep-ph/0411402

The Platinum Channel  $P_{\mu e}$



# Killing the Clones

- Combining LBL data with data from other experiments:

Adding reactor antineutrino data

Huber, Lindner, Schwetz, Winter, hep-ph/0303232

Adding atmospheric neutrino data

Huber, Maltoni, Schwetz, hep-ph/0501037

Campagne, Maltoni, Mezzetto, Schwetz, hep-ph/0603172



# Killing the Clones at The Magic Baseline

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# The Magic baseline

$$\begin{aligned} P_{e\mu} &\simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \frac{\sin^2[(1 - \hat{A})\Delta]}{(1 - \hat{A})^2} \\ &\pm \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \sin \delta_{CP} \sin(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 - \hat{A})\Delta]}{(1 - \hat{A})} \\ &+ \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos \delta_{CP} \cos \Delta \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1 - \hat{A})\Delta]}{(1 - \hat{A})} \\ &+ \alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2} \end{aligned}$$



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$$+ \alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}$$

- If  $\sin(\hat{A}\Delta) \simeq 0$



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- “Clean” measurement of  $\theta_{13}$  and  $\text{sgn}(\Delta m_{31}^2)$



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$$\sin(\hat{A}\Delta) \simeq 0$$

$\Rightarrow$

$$L_{magic} \simeq 7690 \text{ km}$$

Barger, Marfatia, Whisnant, hep-ph/0112119

Huber, Winter, hep-ph/0301257

Smirnov, hep-ph/0610198



# Near-Resonant Matter Effects

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- Large Distance  $\Rightarrow$  Large Matter effects



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- Resonance energy

$$E_{res} = \frac{|\Delta m_{31}^2| \cos 2\theta_{13}}{2\sqrt{2}G_F N_e}$$



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- For  $\Delta m_{31}^2 = 2.5 \times 10^{-3} \text{ eV}^2$ ,  $\sin^2 2\theta_{13} = 0.1$  and the PREM profile  $\rho_{av} = 4.13 \text{ gm/cc}$ ,  $E_{res} \simeq 7.5 \text{ GeV}$



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- Maximal oscillations when  $\sin^2 2\theta_{13}^m \simeq 1$  and  $\sin^2 \left( \frac{(\Delta m_{31}^2)^m L}{4E} \right) \simeq 1$  simultaneously

Gandhi et al, hep-ph/0408361



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- Large Distance  $\Rightarrow$  Large Matter effects
- Resonance energy

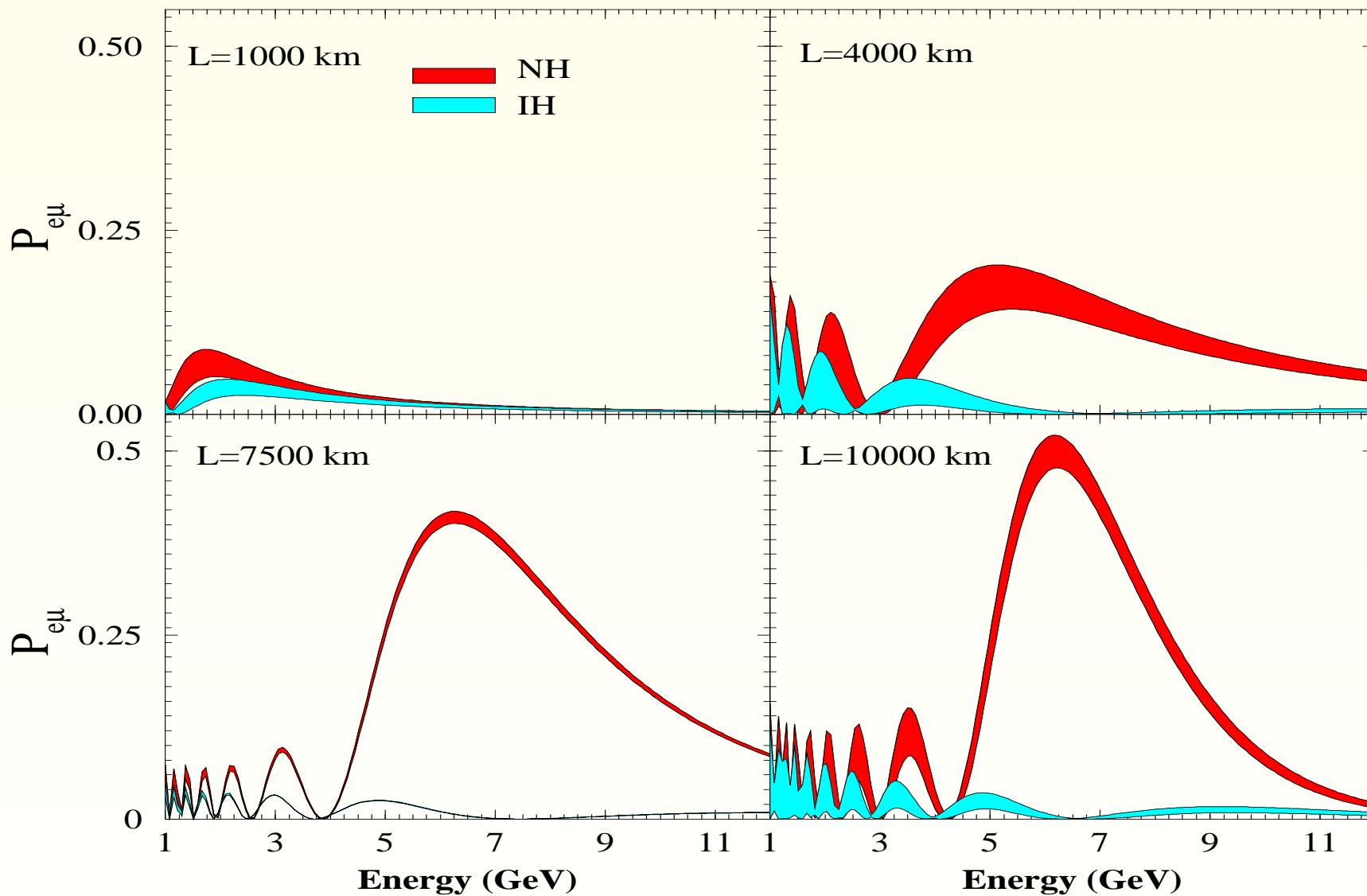
$$E_{res} = \frac{|\Delta m_{31}^2| \cos 2\theta_{13}}{2\sqrt{2}G_F N_e}$$

- For  $\Delta m_{31}^2 = 2.5 \times 10^{-3} \text{ eV}^2$ ,  $\sin^2 2\theta_{13} = 0.1$  and the PREM profile  $\rho_{av} = 4.13 \text{ gm/cc}$ ,  $E_{res} \simeq 7.5 \text{ GeV}$
- Maximal oscillations when  $\sin^2 2\theta_{13}^m \simeq 1$  and  $\sin^2 \left( \frac{(\Delta m_{31}^2)^m L}{4E} \right) \simeq 1$  simultaneously
- At the magic baseline, largest oscillations come when  $E \simeq 6 \text{ GeV}$

Gandhi et al, hep-ph/0408361



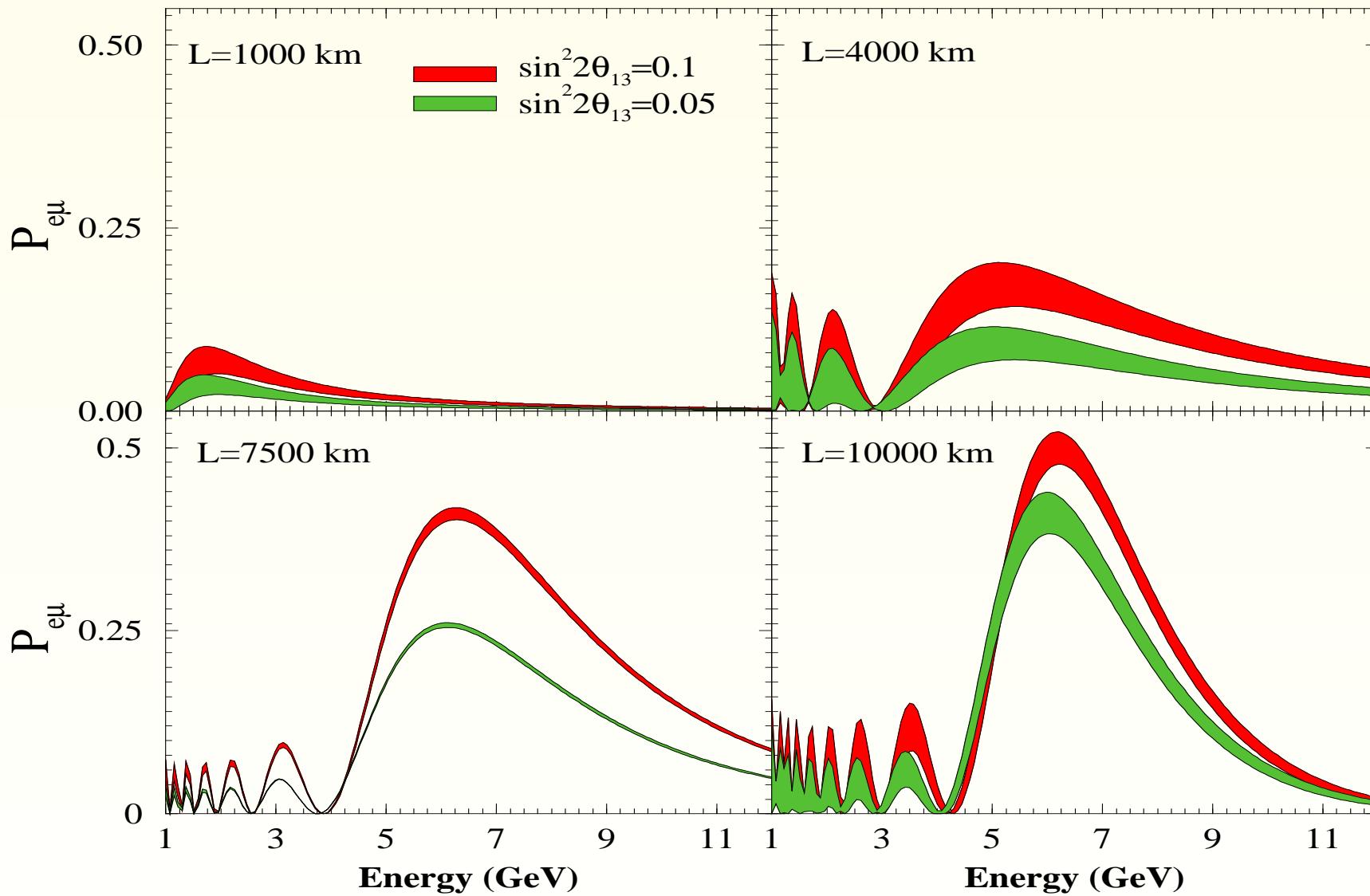
# The Probability



Agarwalla, S.C., Raychaudhuri, hep-ph/0610333



# The Probability



Agarwalla, S.C., Raychaudhuri, hep-ph/0610333



# Conclusions (Comparison of different setups)

	$\gamma$	L(km)	Detector	$T_\nu/T_{\bar{\nu}}$	$\sin^2 2\theta_{13}$	$sgn(\Delta m_{31}^2)$	Max CPV
NF@3000		3000	50 (MI)	4/4	$2.5 \times 10^{-3}$	$(0.8 - 10) \times 10^{-3}$	$7 \times 10^{-5}$
	NF@7500	7500	50 (MI)	4/4	$2 \times 10^{-4}$	$2 \times 10^{-4}$	No sens
CERN-INO	350	7152	50 (MI)	10	$2.1 \times 10^{-3}$	$1.1 \times 10^{-2}$	No sens
	500	7152	50 (MI)	10	$8.4 \times 10^{-4}$	$8.5 \times 10^{-3}$	No sens
hep-ph/ 0603172	100/100	130	440 (WC)	10/10	$5 \times 10^{-3}$ (W) $3 \times 10^{-4}$ (B)	$2.5 \times 10^{-3}$ +SPL+ATM	$2 \times 10^{-4}$
hep-ph/ 0506237	200/200	520	500 (WC)	8/8	$1.5 \times 10^{-3}$	$(0.7 - 2) \times 10^{-2}$	$2 \times 10^{-4}$
	500/500	650	50 (TASD)	8/8	$1.5 \times 10^{-3}$	$(0.6 - 4.5) \times 10^{-2}$	$1 \times 10^{-4}$
	1000/1000	1300	50 (TASD)	8/8	$4 \times 10^{-4}$	$(1 - 7) \times 10^{-3}$	$7 \times 10^{-5}$
hep-ph/ 0312068	100/60	130	400 (WC)	10(S)	Not	No Sens	$1 \times 10^{-3}$
	580/350	732	400 (WC)	10(S)	Given	$2 \times 10^{-2}$	$2 \times 10^{-4}$
	2500/1500	3000	40 (MI)	10(S)		$4 \times 10^{-3}$	$4 \times 10^{-4}$
hep-ph/ 0503021	120/120	130	440 (WC)	10(S)	$5 \times 10^{-3}$	Not	$1 \times 10^{-3}$
	150/150	300	440 (WC)	10(S)	$6 \times 10^{-4}$	Given	$2 \times 10^{-4}$
	350/350	730	440 (WC)	10(S)	$4 \times 10^{-4}$		$1 \times 10^{-4}$



# Backup Slides

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# The CERN-INO Beta-Beam Experiment

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- Proposal to build a large iron detector (ICAL) at INO



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- This is tantalizingly close to the magic baseline



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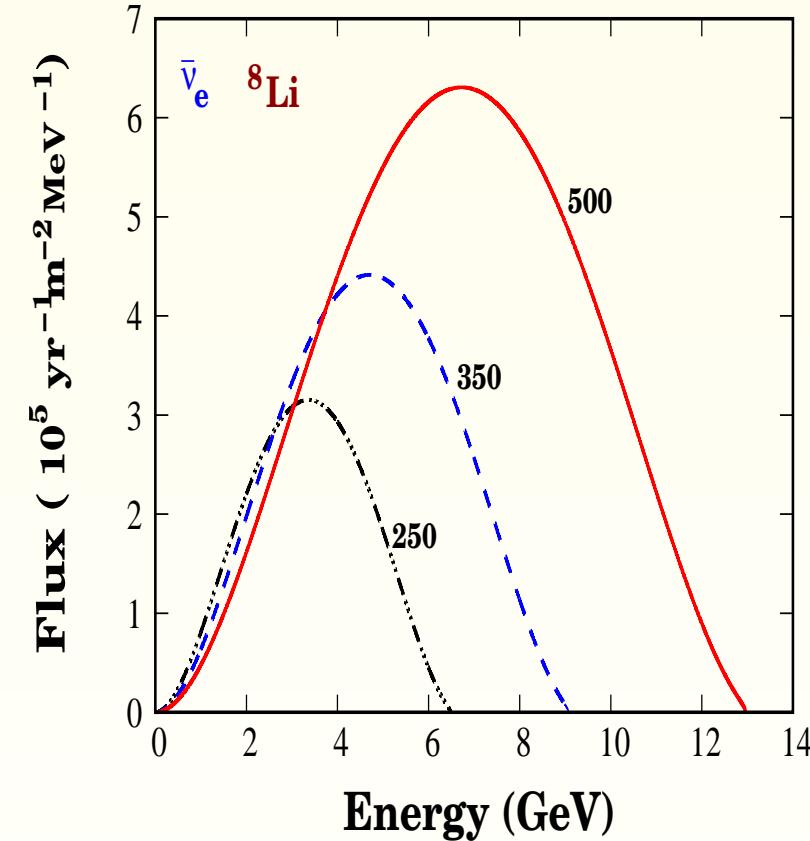
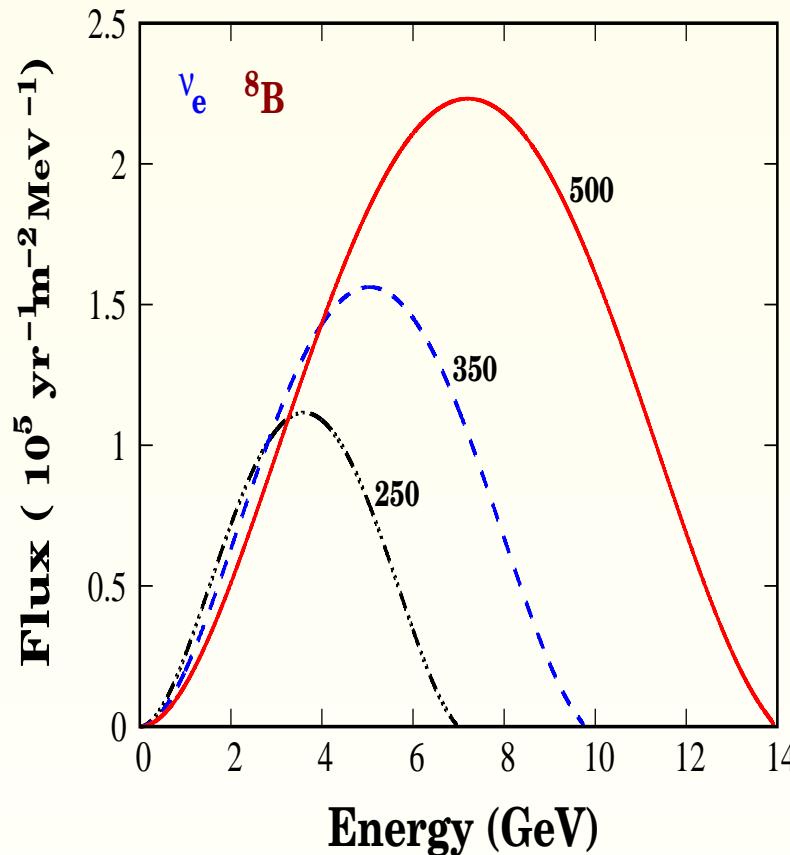


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- Beta beam spectrum depends on the end point energy of the beta unstable ion and Lorentz boost  $\gamma$
- The standard Beta-Beam ions  $^{18}Ne$  and  $^6He$  would require very large gamma
- Alternative ions  $^8B$  and  $^8Li$  have large end-point energy and hence “harder” spectra. Works!!



# The CERN-INO Beta-Beam Experiment

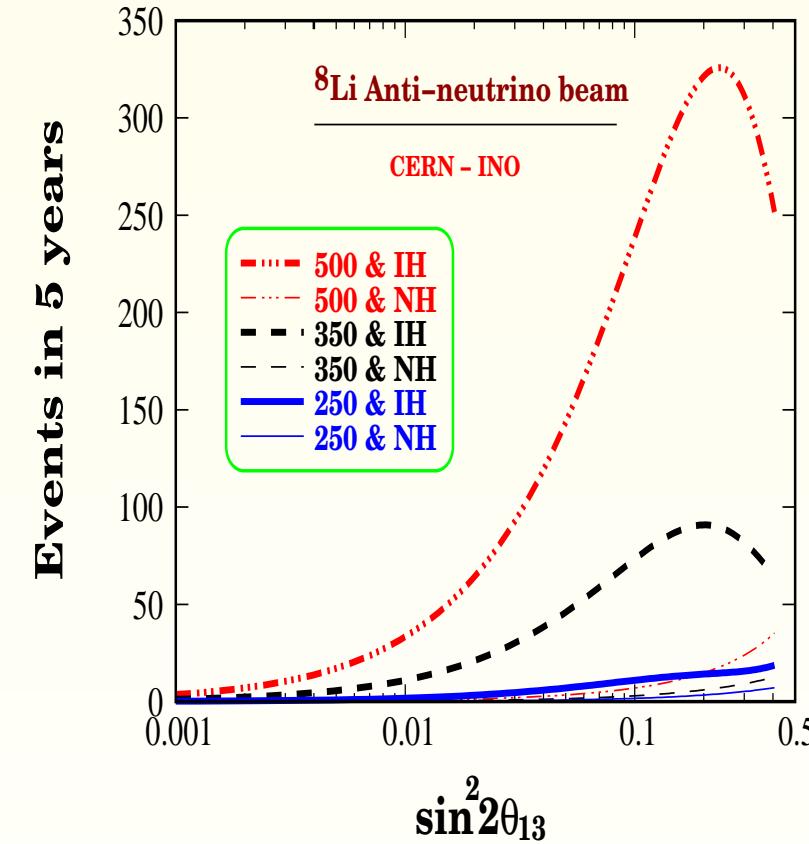
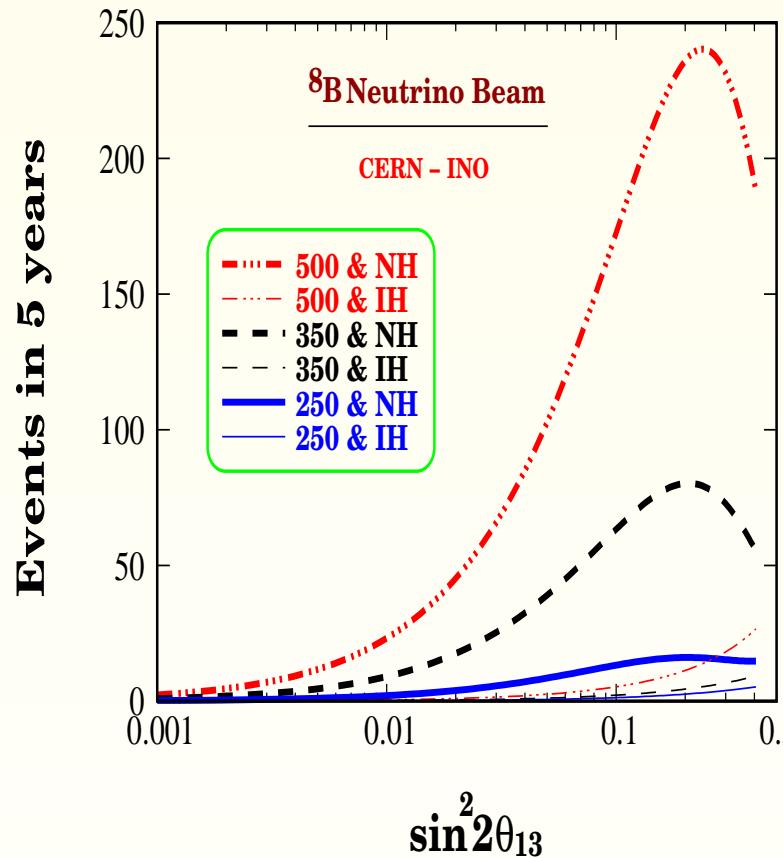


Agarwalla, SC, Raychaudhuri, hep-ph/0610333

- Flux peaks at  $E \simeq 6 \text{ GeV}$  for  $\gamma = 350 - 500$



# The CERN-INO Beta-Beam Experiment



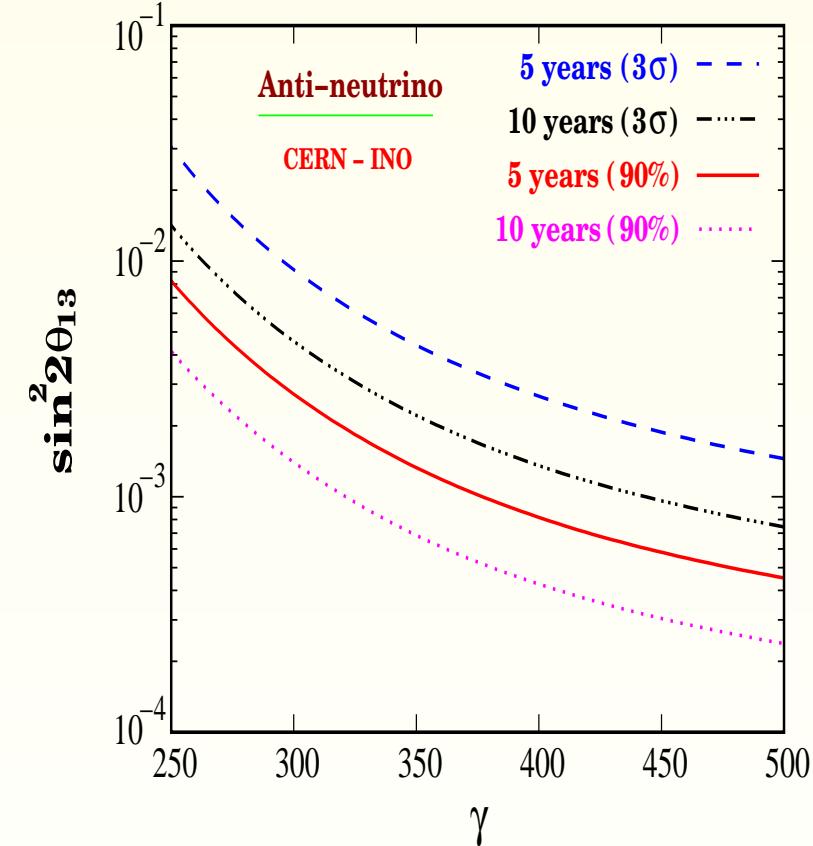
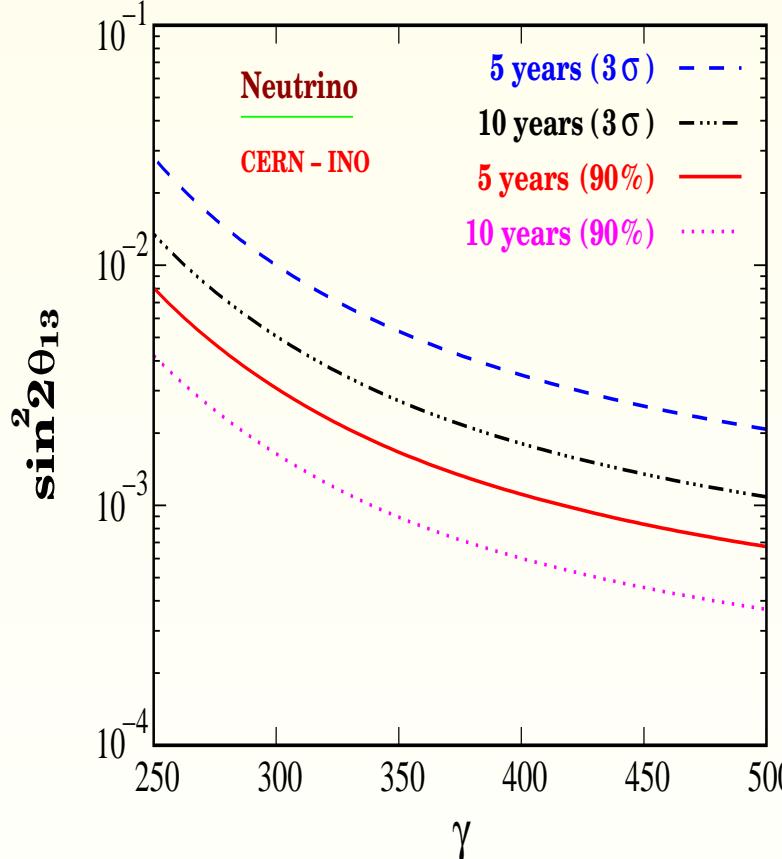
Agarwalla, SC, Raychaudhuri, hep-ph/0610333

- The rate shows a sharp dependence on the hierarchy and  $\theta_{13}$



# The CERN-INO Beta-Beam Experiment

- Sensitivity to  $\theta_{13}$



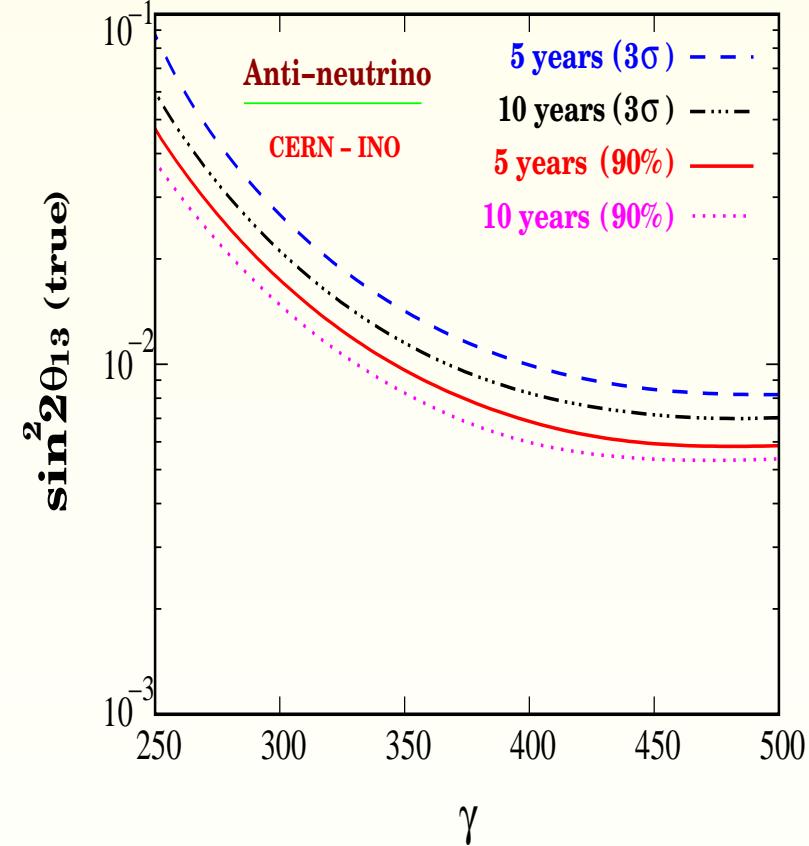
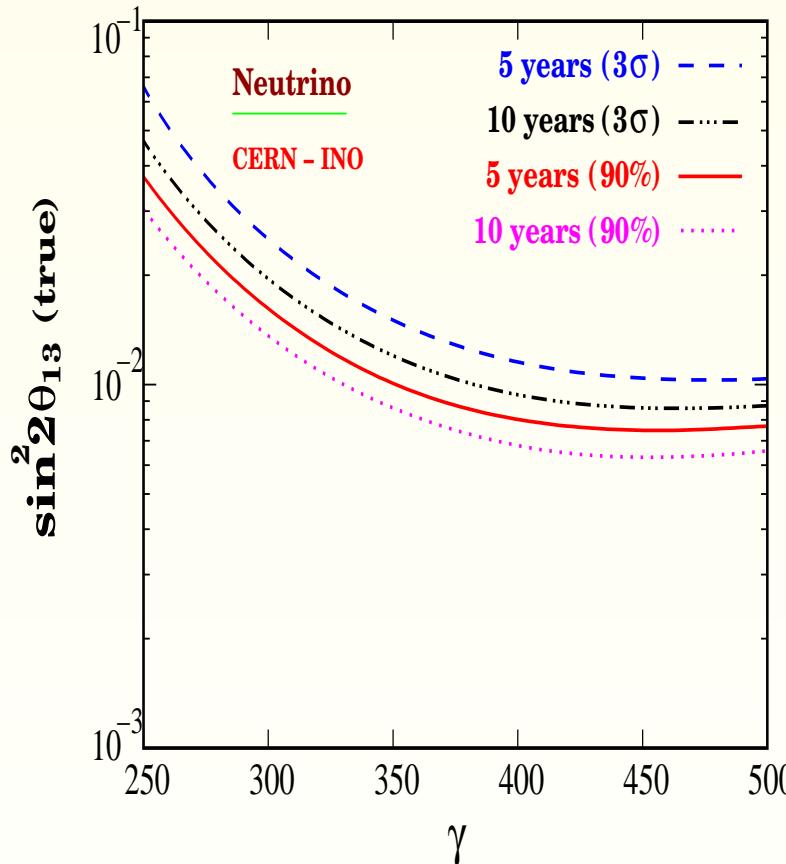
Agarwalla, SC, Raychaudhuri, hep-ph/0610333

- At  $3\sigma$ ,  $\sin^2 2\theta_{13} < 8.5 \times 10^{-4}$  ( $1.5 \times 10^{-3}$ ) with 80% efficiency and 10(5) years data



# The CERN-INO Beta-Beam Experiment

- Sensitivity to  $\text{sgn}(\Delta m_{31}^2)$



Agarwalla, SC, Raychaudhuri, hep-ph/0610333

- At  $3\sigma$ ,  $\sin^2 2\theta_{13} < 8.5 \times 10^{-3} (9.8 \times 10^{-3})$  with 80% efficiency and 10(5) years data



# Large Matter Effects in $P_{ee}$



# Large Matter Effects in $P_{ee}$

$$\lim_{\Delta m_{21}^2 \rightarrow 0} P_{e\mu} = \sin^2 \theta_{23} \sin^2 2\theta_{13}^m \sin^2 \frac{(\Delta m_{31}^2)^m L}{4E}$$

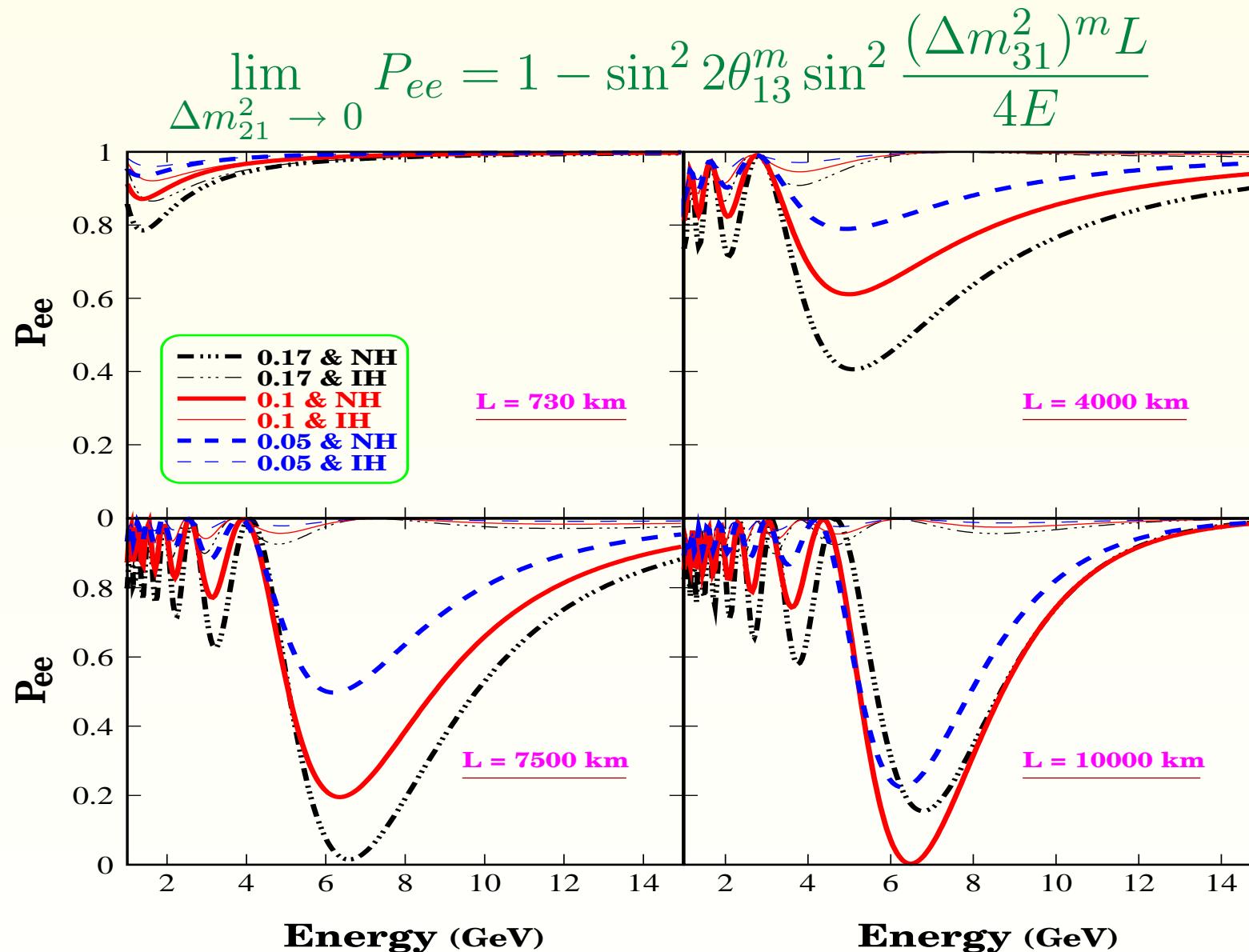


# Large Matter Effects in $P_{ee}$

$$\lim_{\Delta m_{21}^2 \rightarrow 0} P_{e\tau} = \cos^2 \theta_{23} \sin^2 2\theta_{13}^m \sin^2 \frac{(\Delta m_{31}^2)^m L}{4E}$$



# Large Matter Effects in $P_{ee}$

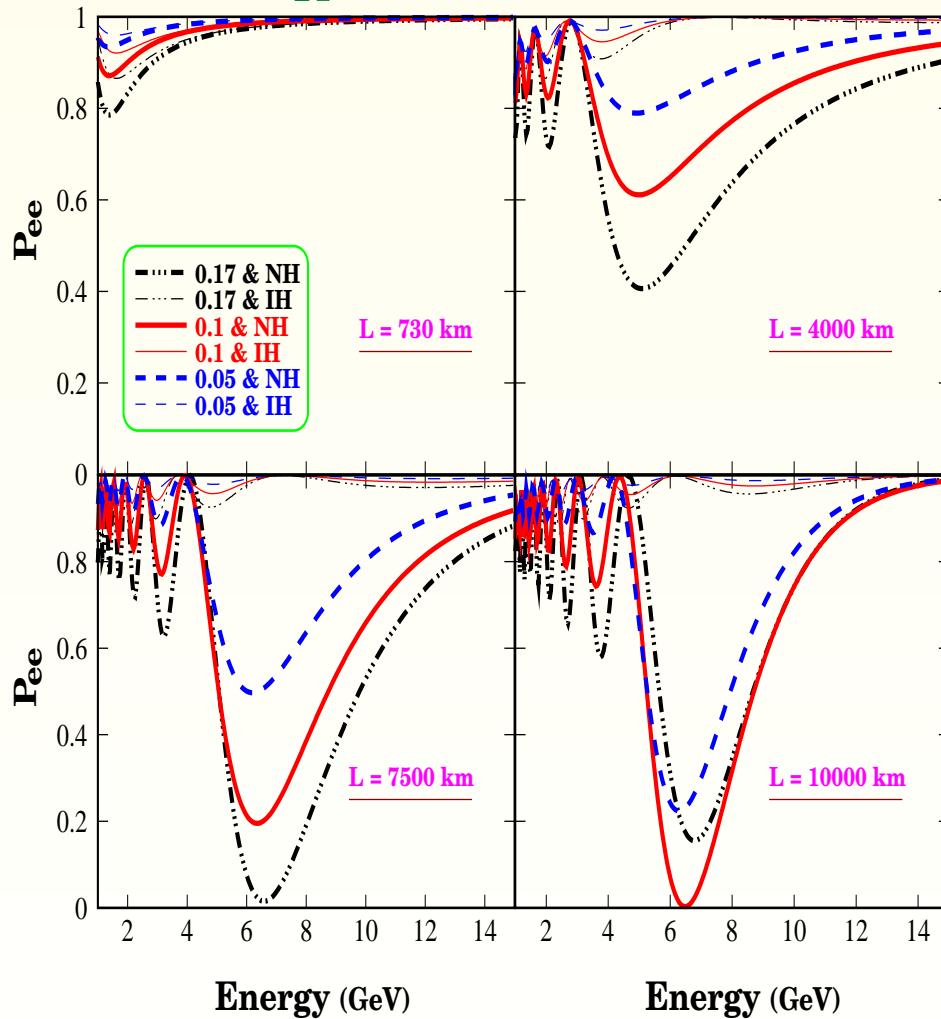


Agarwalla, S.C., Goswami, Raychaudhuri, hep-ph/0611233



# Large Matter Effects in $P_{ee}$

$$\lim_{\Delta m_{21}^2 \rightarrow 0} P_{ee} = 1 - \sin^2 2\theta_{13}^m \sin^2 \frac{(\Delta m_{31}^2)^m L}{4E}$$

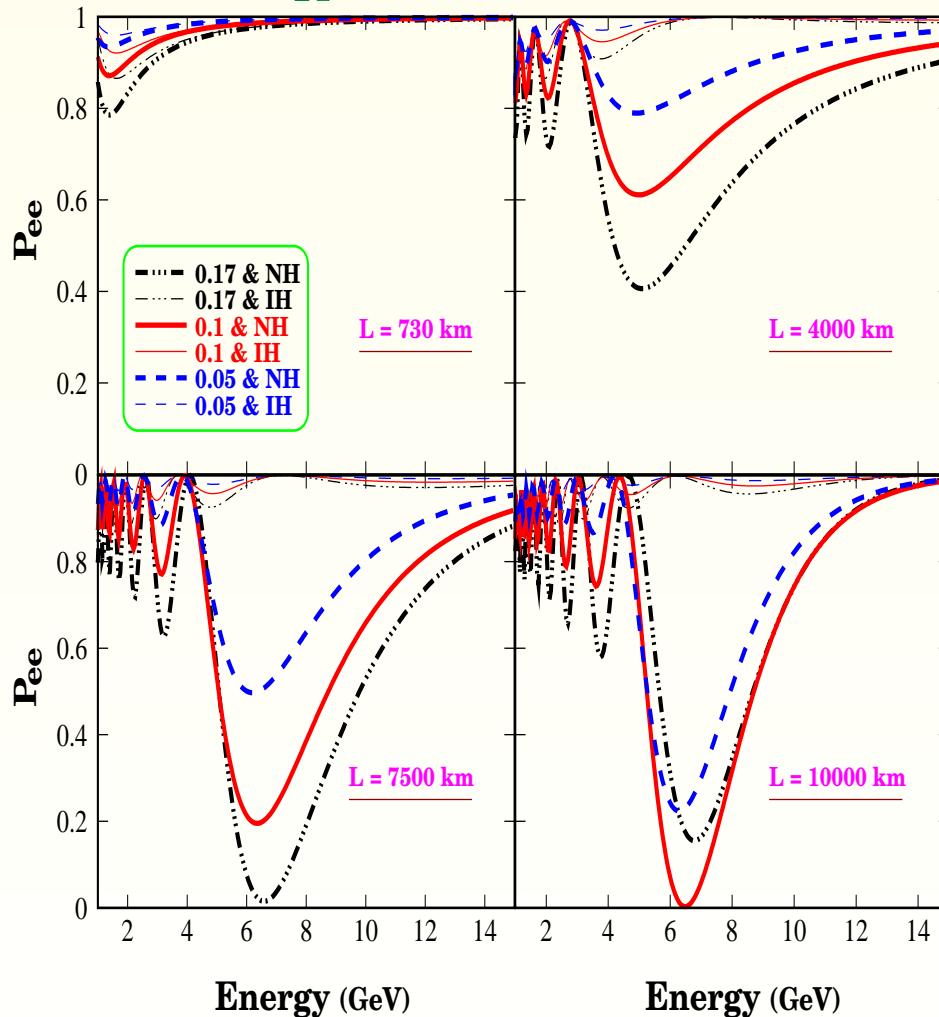


Agarwalla, S.C., Goswami, Raychaudhuri, hep-ph/0611233



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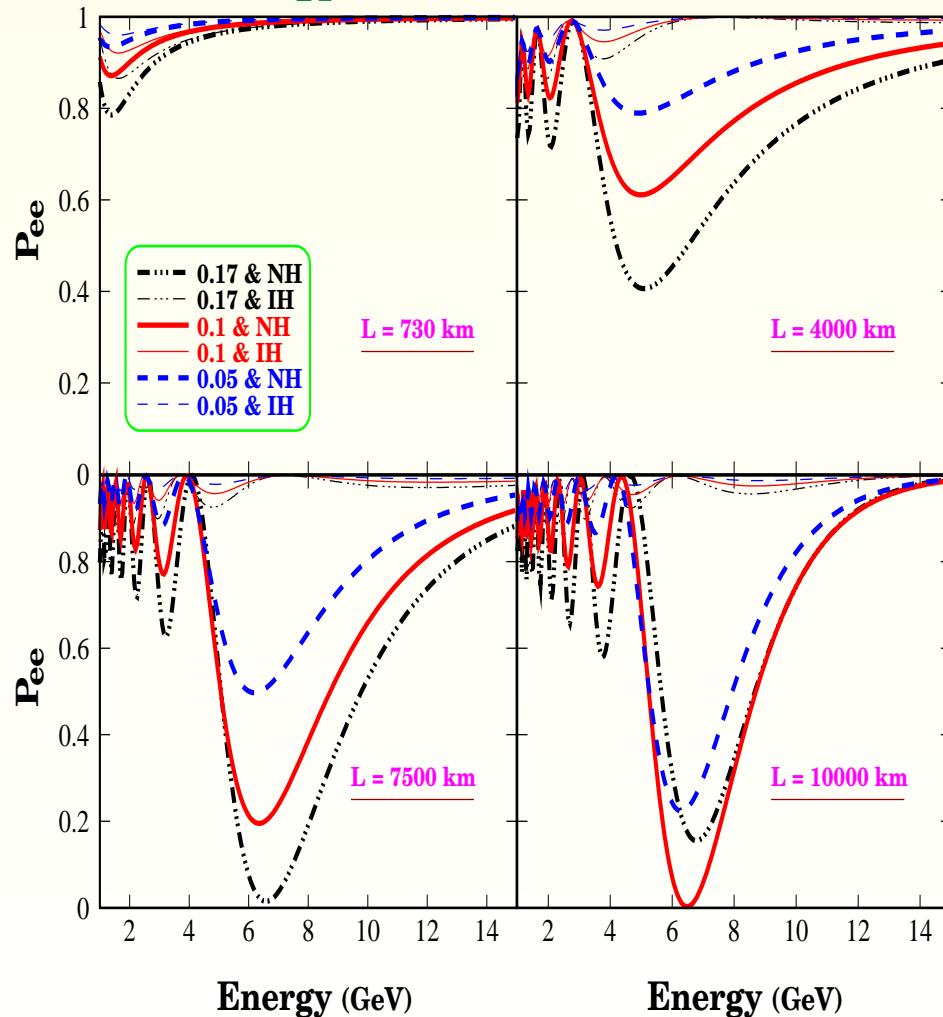
- Matter effects large: almost 2 times
- No dependence on  $\delta_{CP}$  and  $\theta_{23}$
- No parameter degeneracies!!

Agarwalla, S.C., Goswami, Raychaudhuri, hep-ph/0611233



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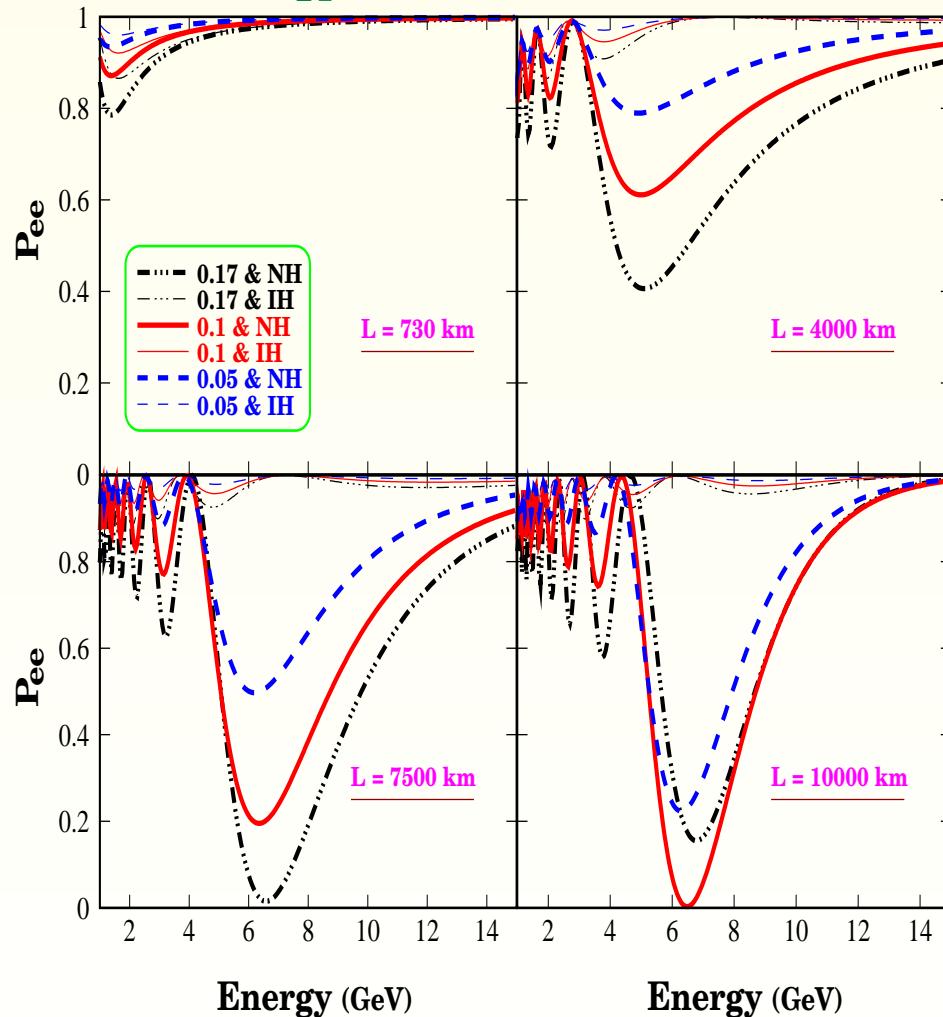
● Pure  $\nu_e/\bar{\nu}_e$ : Beta Beams

Agarwalla, S.C., Goswami, Raychaudhuri, hep-ph/0611233



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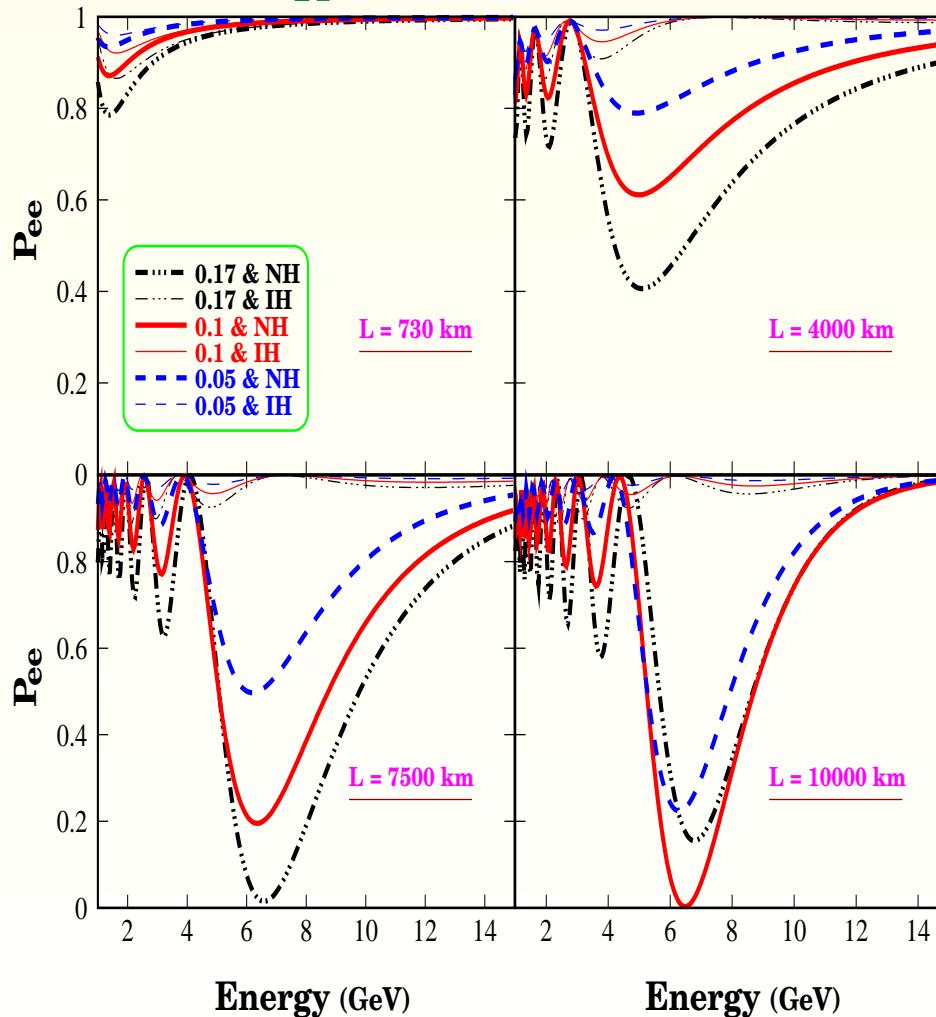
- Pure  $\nu_e/\bar{\nu}_e$ : Beta Beams
- $E \sim 6 \text{ GeV}$ :  ${}^8B$  and  ${}^8Li$

Agarwalla, S.C., Goswami, Raychaudhuri, hep-ph/0611233



# Large Matter Effects in $P_{ee}$

$$\lim_{\Delta m_{21}^2 \rightarrow 0} P_{ee} = 1 - \sin^2 2\theta_{13}^m \sin^2 \frac{(\Delta m_{31}^2)^m L}{4E}$$



- Pure  $\nu_e/\bar{\nu}_e$ : Beta Beams
- $E \sim 6$  GeV:  ${}^8B$  and  ${}^8Li$
- Very long baselines:
  - CERN-UNO:  $L = 7000 - 8600$  km
  - FNAL-MEMPHYS:  $L = 7313$  km
  - FNAL-HK:  $L = 10184$  km
  - CERN-HK:  $L = 9647$  km

Agarwalla, S.C., Goswami, Raychaudhuri, hep-ph/0611233



# CP Discovery Reach with Beta-Beams ( $P_{e\mu}$ )

- At magic baseline CP sensitivity is smothered



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- Must move away from the magic baseline



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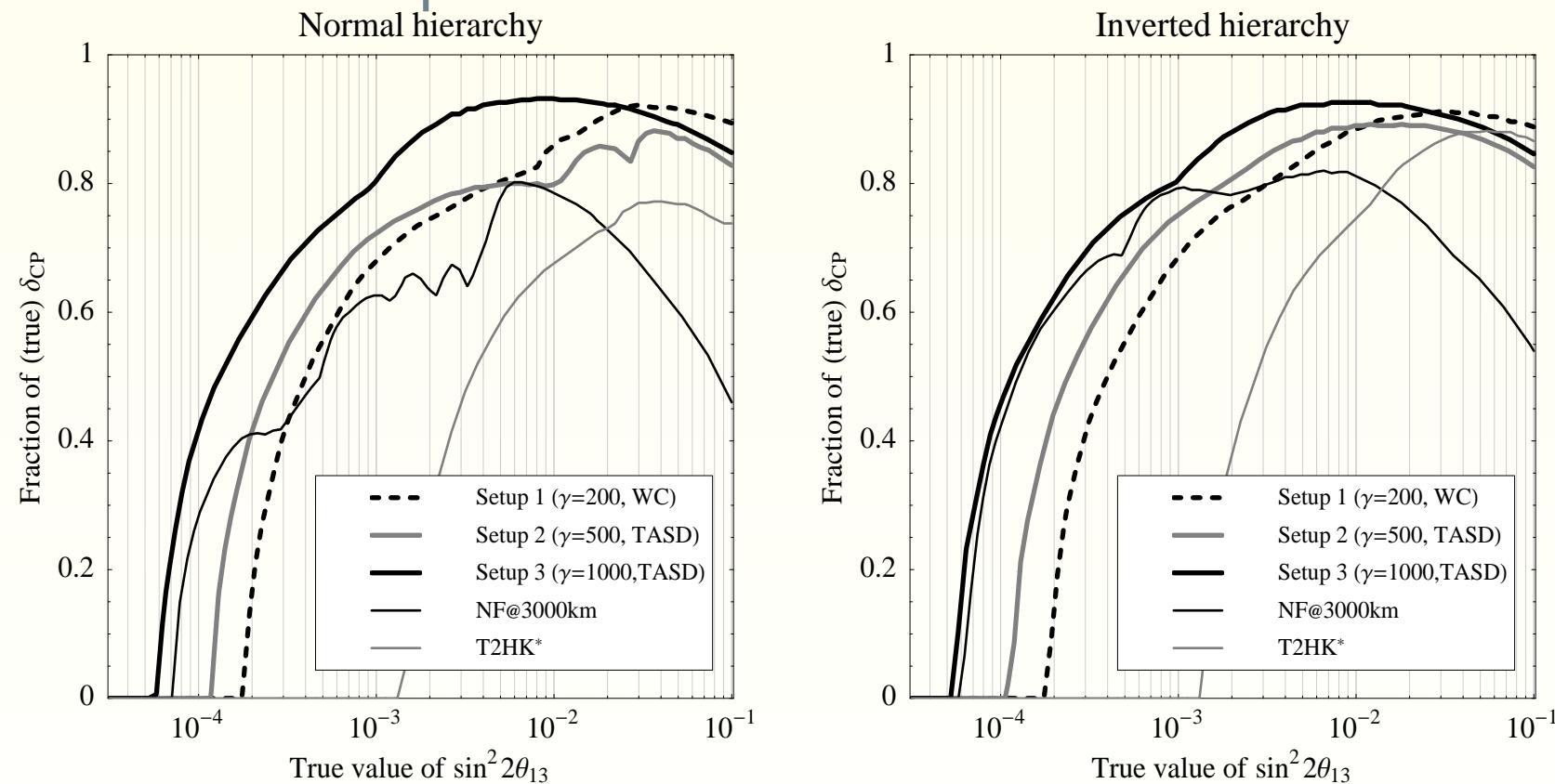
#	1	2	3
type	WC	TASD	TASD
$m$ [kt]	500	50	50
$\gamma$	200	500	1000
$L$ [km]	520	650	1000
$\nu$ signal	1983	2807	7416
$\nu$ background	105	31	95

The following results are taken from  
PH, M. Lindner, M. Rolinec, W. Winter,  
[hep-ph/0506237](https://arxiv.org/abs/hep-ph/0506237).



# CP Discovery Reach with Beta-Beams ( $P_{e\mu}$ )

- At magic baseline CP sensitivity is smothered
- Must move away from the magic baseline
- Which set-up could be best?

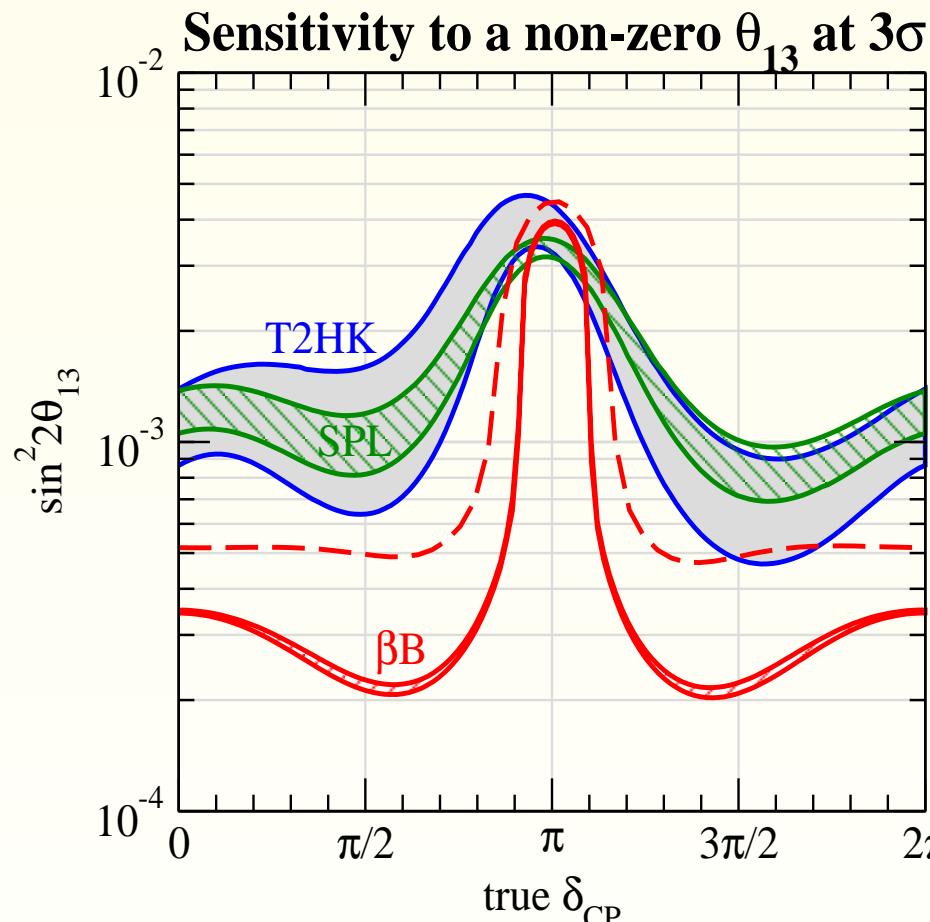


Huber *et al.*, hep-ph/0506237

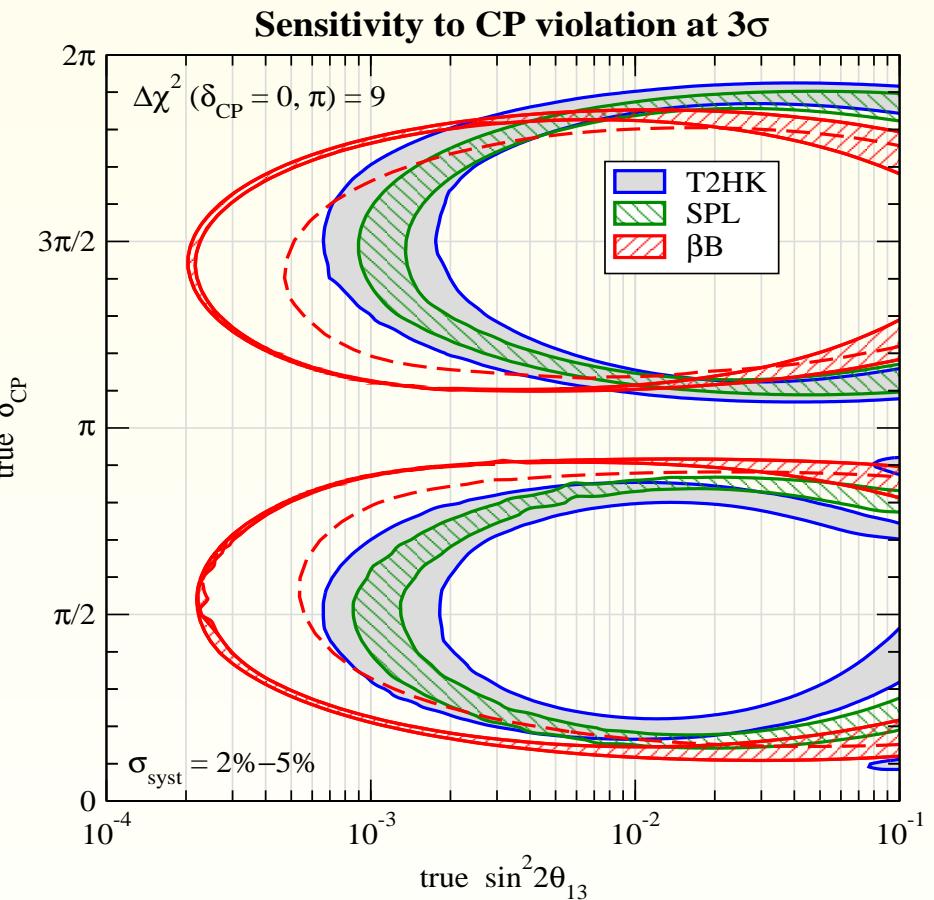
- Beta-Beams CP sensitivity is similar to NuFact



# Physics Potential of CERN-MEMPHYS



Campagne *et al.*, hep-ph/0603172

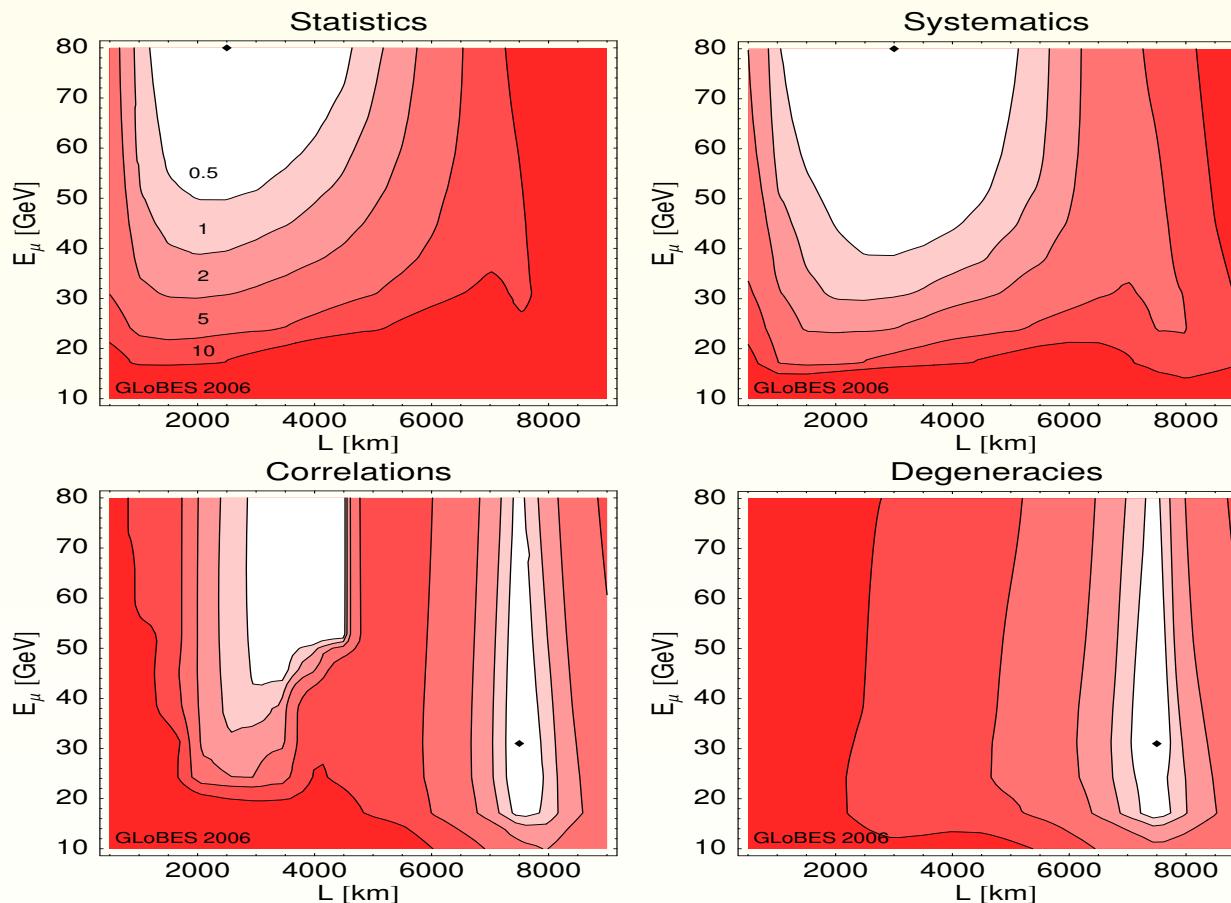


- $\gamma = 100, L = 130$  km, 440 kton WC detector
- Combined  $\beta$ B and SPL in only  $\nu$  mode takes 5 yrs only
- The megaton water detector will also collect atm data  
– helps in determining  $sgn(\Delta m_{31}^2)$  and octant of  $\theta_{23}$



# Physics Reach of Neutrino Factory

- Sensitivity to  $\sin^2 2\theta_{13} \lesssim 5.0 \times 10^{-4}$  ( $5\sigma$ )



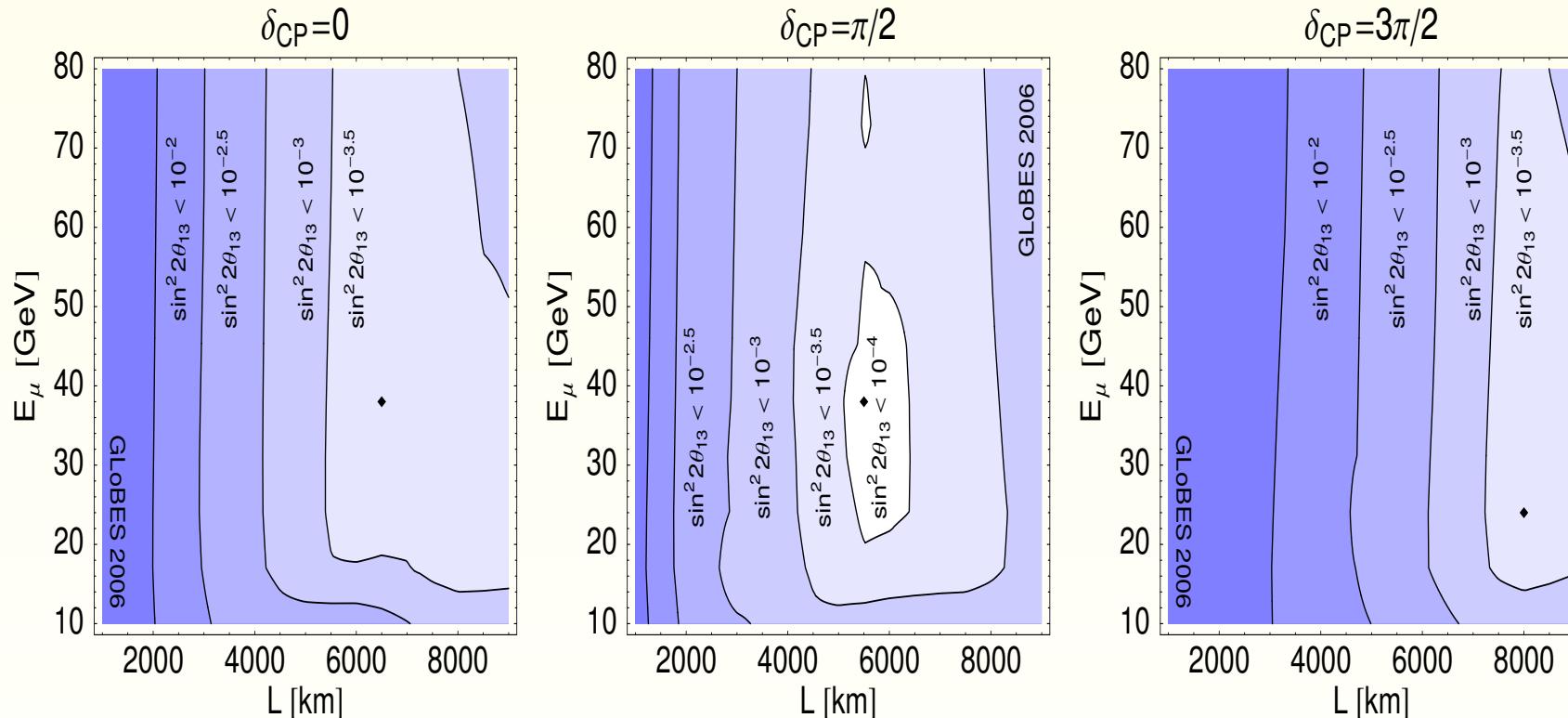
Huber *et al.*, hep-ph/0606199

- Best around  $E_\mu = 30$  GeV and close to magic baseline



# Physics Reach of Neutrino Factory

- Sensitivity to  $\text{sgn}(\Delta m_{31}^2) \gtrsim 1.8 \times 10^{-4}$  ( $3\sigma$ )



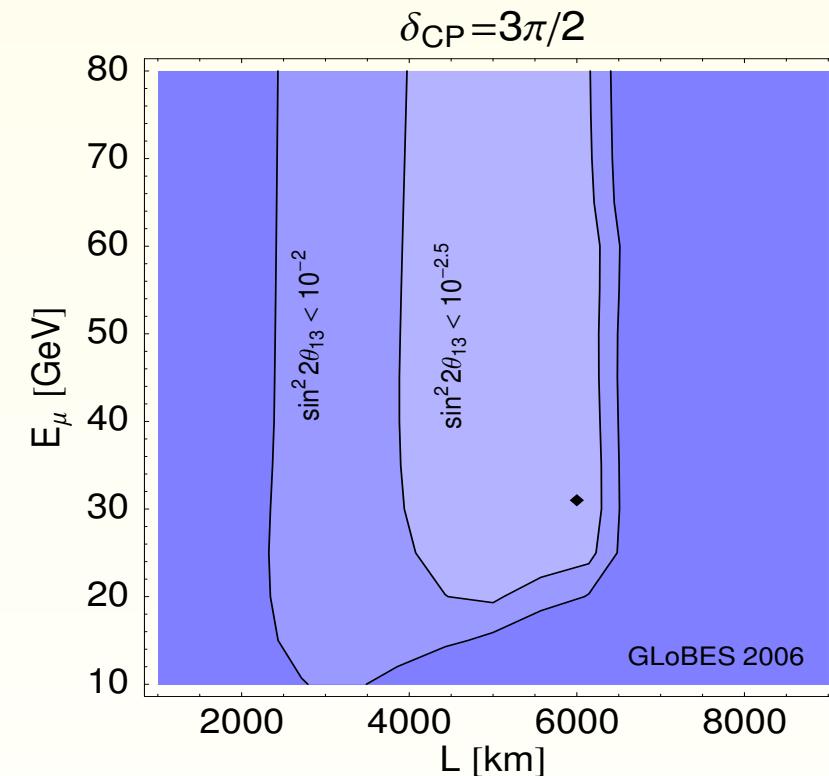
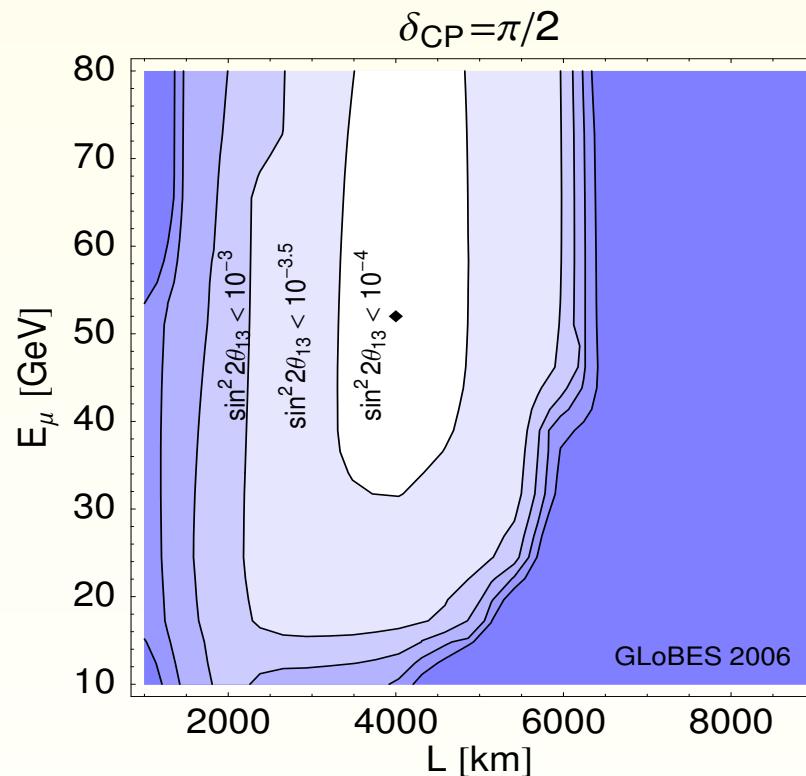
Huber *et al.*, hep-ph/0606199

- Best around  $E_\mu = 20 - 40$  GeV and close to magic baseline



# Physics Reach of Neutrino Factory

- CP Discovery Reach ( $3\sigma$ )



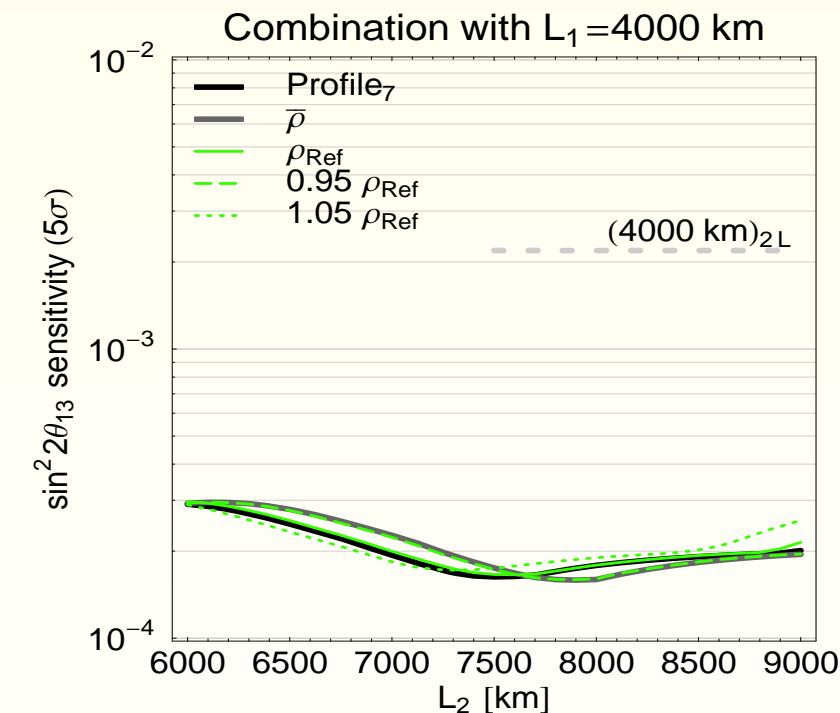
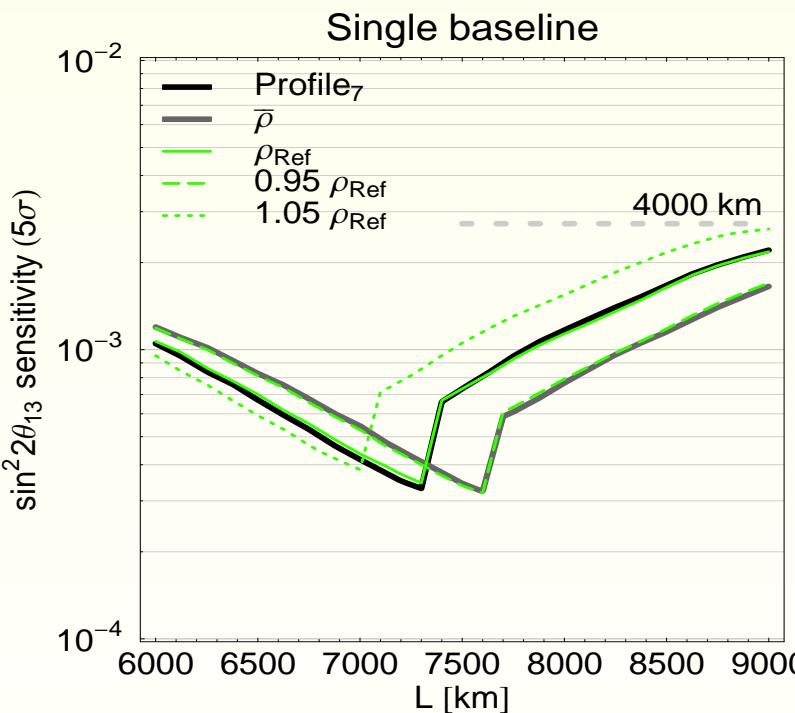
Huber *et al.*, hep-ph/0606199

- Best at around  $L = 4000$  km



# Physics Reach of Neutrino Factory

- Combining baselines for best  $\theta_{13}$  sensitivity



Gandhi, Winter, hep-ph/0612158

- With two detectors, one at  $L = 4000$  km and another farther away, the range of “optimal” baselines widens