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ATLAS CALORIMETER PERFORMANCE FOR CHARGED PION

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The intrinsic performance of the ATLAS barrel and extended barrel calorimeters for the measurement of charged pions is presented. Pion energy scans ($E = 20, 50, 200, 400, \text{ and } 1000 \text{ GeV}$) at two pseudo-rapidity points ($\eta = 0.3 \text{ and } 1.3$) and pseudorapidity scans ($-0.2 < \eta < 1.8$) with pions of constant transverse energy ($E_T = 20 \text{ and } 50 \text{ GeV}$) are analysed. A simple approach, that accounts in first order for noncompensation and dead material effects, is used for the pion energy reconstruction. The intrinsic performances of the calorimeter are studied: resolution, linearity, effect of dead material, tails in the energy distribution. The effect of electronic noise, cell energy cuts and restricted cone size are investigated.

The investigation has been performed at the Dzhelapov Laboratory of Nuclear Problems, JINR.

Свойства АТЛАС калориметрии при регистрации заряженного пиона

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Представлены результаты исследования свойств калориметрического комплекса эксперимента ATLAS при регистрации заряженного пиона. Проведен анализ данных, полученных методом моделирования, для пионов с энергиями 20, 50, 200, 400, 1000 ГэВ и псевдобыстротами 0,3, 1,3; а также для пионов с постоянной поперечной энергией ($E_T = 20 \text{ и } 50 \text{ ГэВ}$) и псевдобыстротами в интервале от $-0,2$ до $1,8$. Для реконструкции энергии пионов был применен простейший метод, в котором учитываются некомпенсированность калориметров и энергетические потери в мертвой материи. Были изучены внутренние свойства АТЛАС калориметрии: разрешение, линейность, эффект мертвой материи, хвосты энергетического распределения. Исследовано влияние на точность реконструкции энергии пиона электронного шума, критериев обрезания по величине сигнала в ячейке калориметра и размеру конуса.

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1. ATLAS CALORIMETRY

The ATLAS calorimetry consists of an electromagnetic (EM) calorimeter covering the pseudorapidity region $|\eta| < 3.2$, a hadronic barrel calorimeter covering $|\eta| < 1.7$, hadronic end-cap calorimeters covering $1.5 < |\eta| < 3.2$, and forward calorimeters covering $3.1 < |\eta| < 4.9$. A view of the ATLAS calorimeters [1] is presented in Fig. 1.

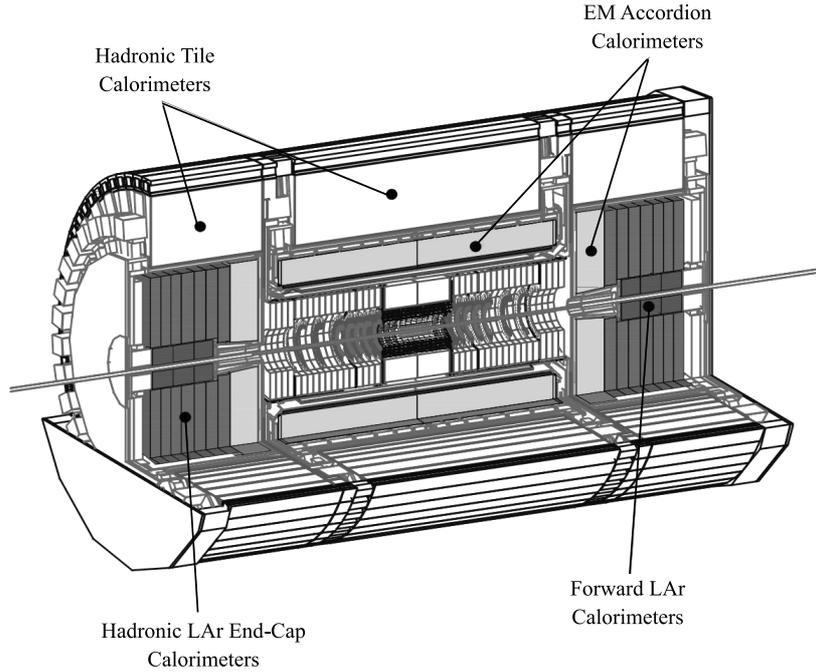


Fig. 1. Three-dimensional cutaway view of the ATLAS calorimeters

The EM calorimeter is a lead/liquid-argon (LAr) detector with accordion geometry [2]. Over the pseudorapidity range $|\eta| < 1.8$, it is preceded by a presampler detector, installed immediately behind the cryostat cold wall, and used to correct for the energy lost in the material upstream of the calorimeter.

The hadronic calorimetry of ATLAS, presented in Figure 1, consists of three main devices. In the barrel region ($|\eta| < 1.7$) there is the scintillating Tile Calorimeter [3]. The Hadronic End-Cap LAr Calorimeter (HEC) extends up to $|\eta| = 3.2$. The range $3.1 < |\eta| < 4.9$ is covered by the high density Forward Calorimeter (FCAL). Up to $|\eta| = 2.5$ the basic granularity of the hadron calorimeters is $\Delta\eta \times \Delta\phi = 0.1 \times 0.1$. This region is used for precise measurements of the energy and angles of jets. In the region $|\eta| > 2.5$, the basic granularity is approximately $\Delta\eta \times \Delta\phi = 0.2 \times 0.2$.

A more detailed description of all ATLAS calorimeters is given in the Calorimeter TDRs ([1], [2] and [3]).

The performance of the barrel and extended barrel sections of the ATLAS hadronic calorimeter for the measurement of charged pion energy is studied. The intrinsic energy resolution, the effects of dead material, electronic noise and limited cone size are discussed.

2. ENERGY RESOLUTION

In the barrel region, the response of the calorimeter was studied at two pseudorapidity values: $\eta = 0.3$ (central barrel) and $\eta = 1.3$ (extended barrel). First, the energy sampled in the different calorimeter compartments is converted to the total deposited energy using the electromagnetic energy scale (EM scale). The intrinsic performance of the calorimeter is studied: the energy considered is not restricted to a cone and electronic noise is not added. These effects are discussed later in Section 4. The algorithm to reconstruct the pion energy is similar to the «Benchmark Method» used to analyse the combined LAr-Tile test beam data [4]

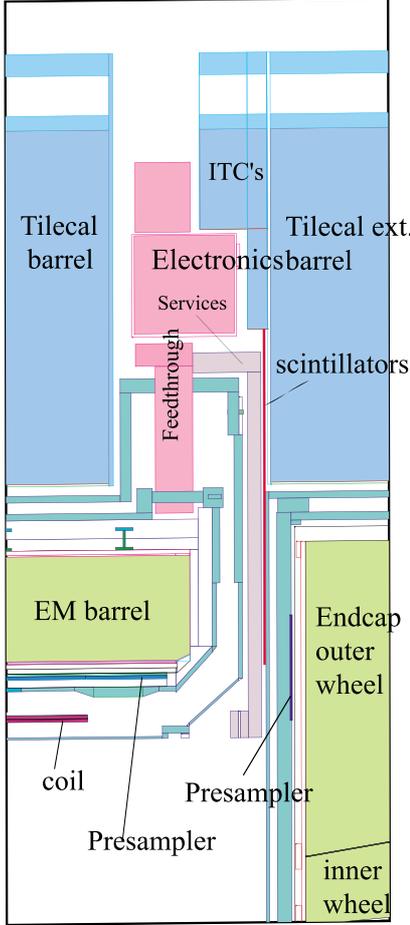


Fig. 2. Longitudinal view of a quadrant of the vertical gap between the barrel and extended barrel Hadronic Tile Calorimeter

$$E_{\text{rec}} = \alpha \cdot E_{\text{had}} + \beta \cdot E_{\text{em}} + \gamma E_{\text{em}}^2 + \delta \cdot \sqrt{E_{\text{had}1} \cdot E_{\text{em}3}} + \kappa \cdot E_{\text{ITC}} + \lambda \cdot E_{\text{scint}} \quad (1)$$

The coefficients α and β take into account the different response of the EM and Hadronic Calorimeters to the pion energy. The quadratic term γE_{em}^2 provides an additional first order correction for non-compensation (the coefficient is negative, it suppresses the signal for events with a large fraction of electromagnetic energy). The term $\delta \sqrt{E_{\text{had}1} \cdot E_{\text{em}3}}$ estimates the energy loss in the cryostat wall separating the LAr and Tile Calorimeters. In the central barrel, the energy is taken from the geometric mean of the energies in the last compartment of the LAr EM barrel ($E_{\text{em}3}$) and the first compartment of the Tile barrel calorimeter ($E_{\text{had}1}$); whereas in the extended barrel the energy is taken from the geometric mean of the energies in the outer wheel of the EM end-cap and the first compartment of the Tile extended barrel calorimeter. The term κE_{ITC} corrects for the energy loss in the dead material in the vertical gap between the Tile central and extended barrels. It is sampled by the two Intermediate Tile Calorimeter (ITC) modules (see Fig. 2.). The last term λE_{scint} corrects for the energy loss in the barrel and end-cap vertical cryostat walls (see Figure 2).

The response and the energy resolution for pions in the energy range from $E = 20$ GeV to 1 TeV at $\eta = 0.3$ and 1.3 are shown in Figs. 3 and 4. The open crosses show the results when the coefficients of Equation (1) are independent of energy. With the simple «Benchmark Method», the effect of noncompensation is not fully corrected for and there is a residual nonlinearity of the pion response of the order of 4 – 5% between 20 GeV and 1 TeV. The test beam data show a 10% residual nonlinearity between 20 and 300 GeV when using the

same reconstruction method [4], reflecting the fact that G-CALOR [5] predicts a lower degree of noncompensation and may not describe correctly the energy dependence of the fraction of electromagnetic energy produced in the pion interaction. The energy dependence of the resolution is fitted with the two-term formula

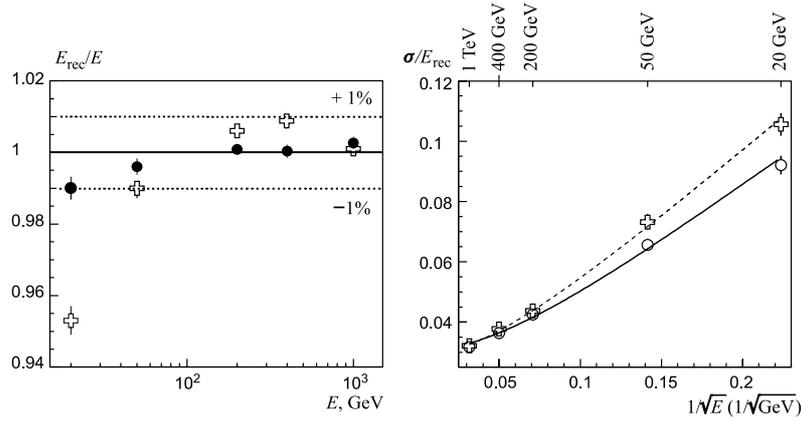


Fig. 3. Pion energy scan in the central barrel ($\eta = 0.3$). The left plot shows the residual nonlinearity, the right plot shows the energy resolution with the results of the fit with Eq. $\sigma/E = a/\sqrt{E} \oplus b$. Two sets of parameters for the pion energy reconstruction have been used: open crosses — for energy independent parameters; solid dots — for parameters fitted at each energy and pseudorapidity

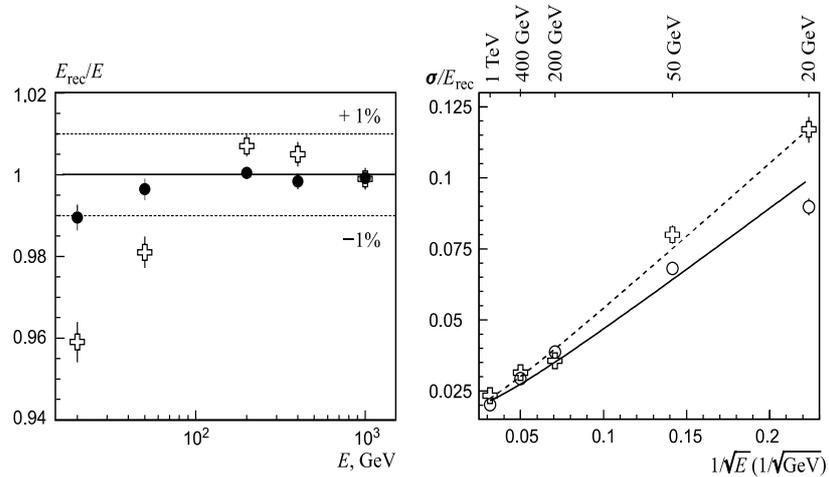


Fig. 4. Pion energy scan in the extended barrel ($\eta = 1.3$). The left plot shows the residual nonlinearity, the right plot shows the energy resolution with the results of the fit with Eq. $\sigma/E = a/\sqrt{E} \oplus b$. Two sets of parameters for the pion energy reconstruction have been used: open crosses — for energy independent parameters; solid dots — for parameters fitted at each energy and pseudorapidity

$$\sigma/E = a/\sqrt{E} \oplus b \quad (2)$$

where the sampling term a is given in $\% \sqrt{\text{GeV}}$ and the constant term b in $\%$. The pion energy resolution for $\eta = 0.3$ is $(40 \pm 1)\% \sqrt{\text{GeV}}/\sqrt{E} \oplus (3.0 \pm 0.1)\%$ and for $\eta = 1.3$ is $(44 \pm 3)\% \sqrt{\text{GeV}}/\sqrt{E} \oplus (1.6 \pm 0.3)\%$. Although the resolutions obtained for low-energy pions are similar in both cases, at high energy there is some longitudinal leakage in the central barrel, yielding a resolution at 1 TeV of 3% instead of the 2% obtained in the extended barrel. When energy dependent parameters are applied (solid dots), the linearity of the response is restored * and the resolution improved.

3. PSEUDORAPIDITY SCAN

A pseudorapidity scan with pions of constant transverse energy $E_T = 20$ and 50 GeV was carried out to check that the linearity of the response can be maintained across the pseudorapidity range covered by the barrel and the extended barrel, and that no significant tail appears in the line shape. The algorithm, characterised by Eq. (1), with energy dependent parameters was applied. The parameters were adjusted independently for the six sets of pion data, each one covering an interval of 0.4 in pseudorapidity.

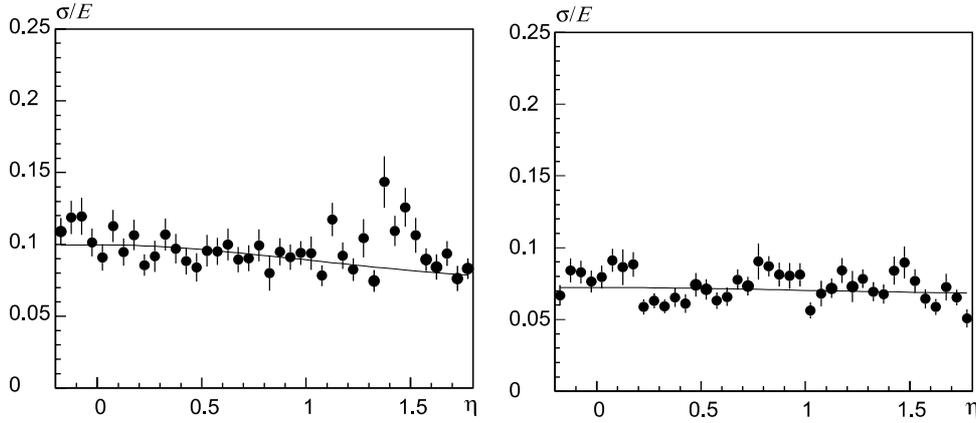


Fig. 5. The dependence of the energy resolution on pseudorapidity for charged pions of constant transverse energy: $E_T = 20$ GeV (left plot) and $E_T = 50$ GeV (right plot). The lines correspond to the energy resolution parametrised using Eq. $\sigma/E = 39\%/\sqrt{E} \oplus 1\%$ for $E_T = 20$ GeV and $\sigma/E = 49\%/\sqrt{E} \oplus 2\%$ for $E_T = 50$ GeV

The energy resolutions obtained for the two scans are shown in Fig. 5. The solid lines show the energy resolution corresponding to Eq. $\sigma/E = 39\%/\sqrt{E} \oplus 1\%$ for $E_T = 20$ GeV and $\sigma/E = 49\%/\sqrt{E} \oplus 2\%$ for $E_T = 50$ GeV. This performance allows one to fulfil the

*The 1% residual nonlinearity at 20 GeV results from the fact that the coefficients were obtained by minimizing the expression $\sum (E_{\text{rec}} - E)^2$ without the addition of a linear term $\sum (E_{\text{rec}} - E)$ with a Lagrange multiplier.

goal for the jet energy resolution of the ATLAS hadronic calorimetry in the region $|\eta| < 3$ of $\sigma/E = 50\%/\sqrt{E} \oplus 3\%$.

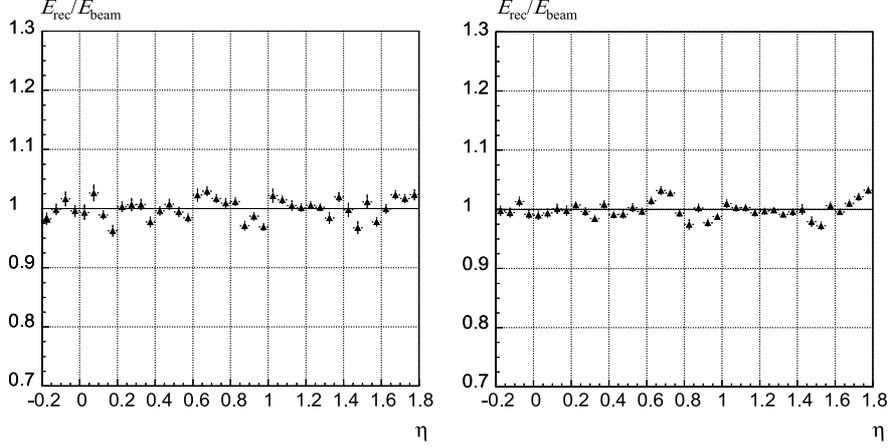


Fig. 6. Single charged pion linearity across η : $E_T = 20$ GeV (left plot) and $E_T = 50$ GeV (right plot). Calibration coefficients are adjusted independently for each 0.4 bin in η but are kept constant within that interval. The figure shows the fitted mean of the pion distribution per bin of 0.05 in η

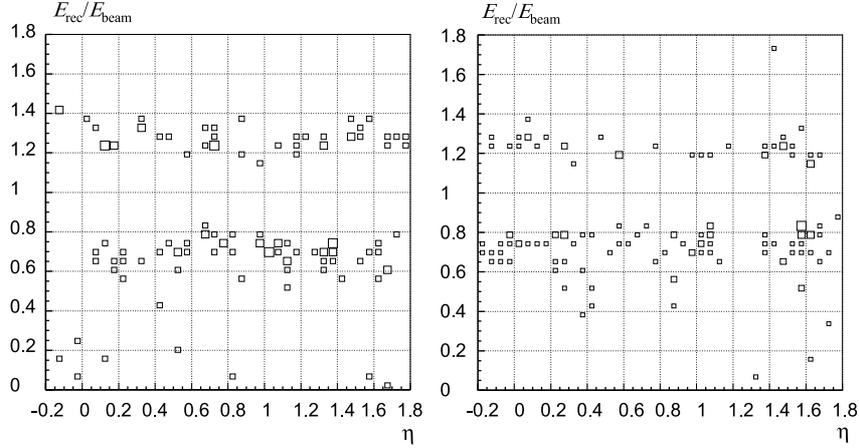


Fig. 7. Events in the tails of the distribution of the reconstructed energy as a function of pseudorapidity for pions of $E_T = 20$ GeV (left plot) and $E_T = 50$ GeV (right plot). Tails are defined as events with reconstructed energies more than three standard deviations away from the mean

In the region of the cracks between the calorimeters, from about $|\eta| = 1.3$ to $|\eta| = 1.5$, where the amount of dead material is the largest, the resolution is somewhat worse.

Figure 6 shows the linearity of the response across η . The fitted mean is plotted for each interval of 0.05 in η . The RMS of the mean is 1.1% for $E_T = 50$ GeV and 2.0% for $E_T = 20$ GeV.

In addition, the tails of the distributions of the reconstructed energy were investigated. Figure 7 shows the events with a pion response more than three standard deviations away from the mean. No significant tails are present: the fraction of events in the tails does not exceed 1 – 2%. A few events out of a total of 5000 events per energy scan, mostly from the sample of pions of $E_T = 20$ GeV, deposit relatively little energy. These correspond to pions decaying to muons before reaching the calorimeter.

4. EFFECTS OF ELECTRONIC NOISE AND CONE SIZE

The results presented so far were obtained without any restriction on the pion reconstruction volume. These results characterise the intrinsic performance of the calorimeters. The presence of electronic noise does not allow integration over a too wide region, therefore the measurement of the pion energy must be restricted to a cone

$$\Delta R = \sqrt{\Delta^2\eta + \Delta^2\phi} . \quad (3)$$

A compromise has to be found between the pion energy lost outside this cone and the noise included inside. The optimum varies as a function of pseudorapidity, since the showers have a width which is characterised by the polar angle whereas the calorimeter cells subtend intervals of constant pseudorapidity. Hence, at higher values of pseudorapidity, the showers extend laterally over more cells.

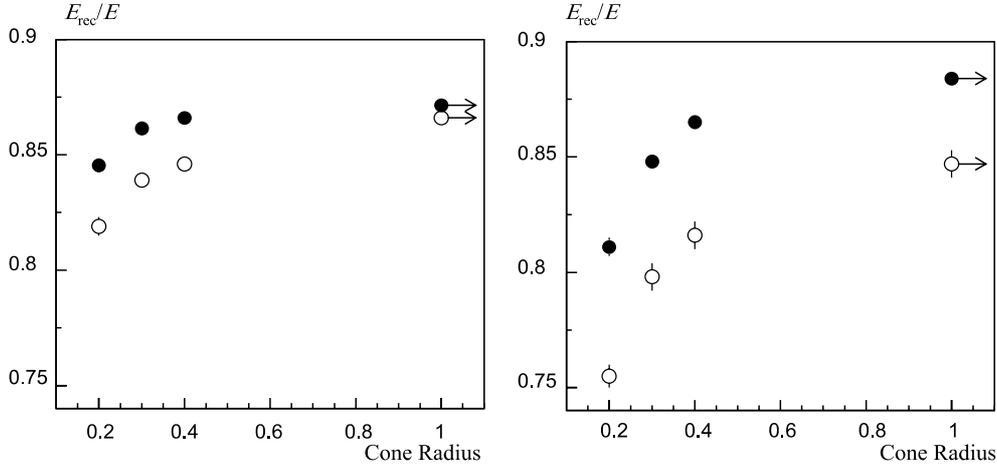


Fig. 8. Energy response for 50 GeV (open circles) and 200 GeV (solid dots) charged pions at $\eta = 0.3$ (left plot) and $\eta = 1.3$ (right plot) as a function of the cone size. The points with arrows correspond to the case without restriction to a cone. Energy and pseudorapidity independent parameters were used for the energy reconstruction

For a cone of $\Delta R = 0.6$ ($\Delta R = 0.3$), the noise is above 3 GeV (1.5 GeV). Digital filtering [6] allows noise suppression (approximately by a factor of 1.6). But even this level of noise is large and is comparable to the intrinsic resolution of the calorimeters for pions

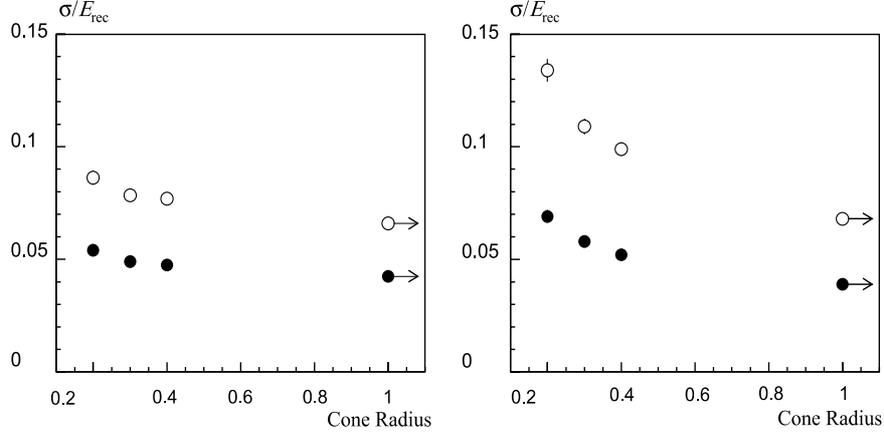


Fig. 9. Energy resolution for 50 GeV (open circles) and 200 GeV (solid dots) charged pions at $\eta = 0.3$ (left plot) and $\eta = 1.3$ (right plot) as a function of the cone size. The points with the arrows correspond to the case without restriction to a cone. Energy and pseudorapidity dependent parameters were used for the energy reconstruction

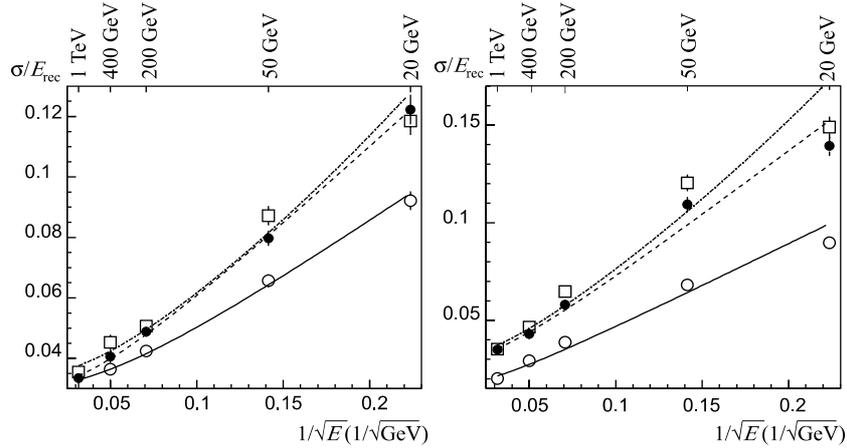


Fig. 10. Energy dependence of the resolution for pions at $\eta = 0.3$ (left plot) and $\eta = 1.3$ (right plot). The open circles show the intrinsic calorimeter resolution obtained with no cut cone restriction and without electronic noise. The solid dots show results obtained for a cone of $\Delta R = 0.3$ and without electronic noise. The open squares show the results obtained for a cone of $\Delta R = 0.3$ and with a 2σ -noise cut (whit electronic noise included). The curves show the results of the fits done with the two-term formula (3) for the first two sets and with the three-term formula (4) for the third set

with energy of a few tens of GeV. A smaller cone of $\Delta R = 0.3$ is preferable from this point of view; after digital filtering, noise can be kept around 1 GeV in the barrel region and below 3 GeV in the pseudorapidity region covered by the extended barrel.

The response and the energy resolution in the barrel region are presented in Figs. 8 and 9 as a function of the cone size used for the pion energy reconstruction. Energy losses outside a

cone noticeably increase with decreasing cone size, especially for 50 GeV pions. The energy resolution also becomes worse, but it is still acceptable for the cone of $\Delta R = 0.3$.

Selecting cells with energy deposition above a certain threshold decreases the noise contribution. A study to optimise the cone size and the noise cut was performed in the barrel region. A 2σ -noise applied to the calorimeter cells within a cone of $\Delta R = 0.3$ leads to the best energy resolution. In Fig. 10, the energy dependency of the resolution is plotted for two pseudorapidities: $\eta = 0.3$ and $\eta = 1.3$. The energy dependence of the resolution can be parametrized by the equation 2 with an additional noise term:

$$\sigma/E = a/\sqrt{E} \oplus b \oplus c/E, \quad (4)$$

where c is given in GeV. The results of the fit with the formula of Eq. 4 are presented in the Table.

Table. Terms of the pion energy resolution fitted with the two-term (2) and the tree-term (4) expression for the $\eta = 0.3$ and 1.3

Central barrel region $\eta = 0.3$			Extended barrel region $\eta = 1.3$		
a (% \sqrt{GeV})	b (%)	c (GeV)	a (% \sqrt{GeV})	b (%)	c (GeV)
No Cone, No Noise					
40 ± 1	3.0 ± 0.1	–	44 ± 3	1.6 ± 0.3	–
Cone $\Delta R = 0.3$, No Noise					
53 ± 2	3.0 ± 0.2	–	67 ± 4	2.9 ± 0.4	–
Cone $\Delta R = 0.3$, Noise with a 2σ -cut					
50 ± 4	3.4 ± 0.3	fixed at 1.0	68 ± 8	3.0 ± 0.7	fixed at 1.5

5. CONCLUSIONS

The response of the barrel and extended barrel region of the ATLAS calorimeter system to single charge pions was investigated using full simulation. Pion energy scans from $E = 20$ GeV to 1000 GeV and pseudorapidity scans with pions of constant transverse energy ($E_T = 20$ and 50 GeV) have been analysed. For the pion energy reconstruction, the «Benchmark approach» was used: it provides a first order correction for noncompensation effects and accounts for the effect of the dead material by using the ITC's and scintillators to sample the energy loss or interpolating between the energy deposited in adjacent calorimeter layers.

Energy and rapidity dependent and independent calibrations have been considered. The best results are obtained with energy and rapidity dependent parameters. The effect of electronic noise has been studied: cone size and cell energy cuts have been optimised. The energy dependence of the resolution can be parameterized as: $(50 \pm 4)\%/\sqrt{E} \oplus (3.4 \pm 0.3)\% \oplus 1.0/E$ at $\eta = 0.3$ and $(68 \pm 8)\%/\sqrt{E} \oplus (3.0 \pm 0.7)\% \oplus 1.5/E$ at $\eta = 1.3$. The larger constant term at $\eta = 0.3$ can be explained by the longitudinal leakage from calorimeters in this region. The resolution, obtained for the pseudorapidity scans, is represented by: $(39 \pm 1)\%/\sqrt{E} \oplus (1 \pm 5)\%$ for $E_T = 20$ GeV, $(49 \pm 9)\%/\sqrt{E} \oplus (2 \pm 6)\%$ for $E_T = 50$ GeV, in the full range, except from about $|\eta| = 1.3$ to $|\eta| = 1.5$, where the resolution is deteriorated by the energy loss in the dead material although no significant tails in the energy spectrum appears.

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